SURFACE CHEMISTRY

Ultrahigh-yield on-surface synthesis and assembly of circumcoronene into a chiral electronic Kagome-honeycomb lattice

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On-surface synthesis has revealed remarkable potential in the fabrication of atomically precise nanographenes. However, surface-assisted synthesis often involves multiple-step cascade reactions with competing pathways, leading to a limited yield of target nanographene products. Here, we devise a strategy for the ultrahigh-yield synthesis of circumcoronene molecules on Cu(111) via surface-assisted intramolecular dehydrogenation of the rationally designed precursor, followed by methyl radical-radical coupling and aromatization. An elegant electrostatic interaction between circumcoronenes and metallic surface drives their self-organization into an extended superlattice, as revealed by bond-resolved scanning probe microscopy measurements. Density functional theory and tight-binding calculations reveal that unique hexagonal zigzag topology of circumcoronenes, along with their periodic electrostatic landscape, confines two-dimensional electron gas in Cu(111) into a chiral electronic Kagome-honeycomb lattice with two emergent electronic flat bands. Our findings open up a new route for the high-yield fabrication of elusive nanographenes with zigzag topologies and their superlattices with possible non-trivial electronic properties.

INTRODUCTION

On-surface synthesis has emerged as a powerful strategy for the bottom-up fabrication of low-dimensional nanographenes (NGs) with atomic precision (1–5). The core concept of the surface-assisted synthetic strategy relies on the rational design of the elegant molecular precursors that subsequently undergo chemical transformation along a certain reaction pathway toward the desired NG product. The versatility of this synthetic strategy has been exemplified by the seminal fabrication of a wealth of open-shell and closed-shell NGs. In addition, the electronic, magnetic, and optical properties of NGs can be precisely tuned via engineering their chemical structure and/or concomitant π -electron topology of the NG (6–14). Therefore, on-surface fabricated NGs with tailored physicochemical properties represent highly prosperous functional materials for next-generation quantum device applications (15–18).

However, on-surface synthetic strategies have long been constrained by a low scalability due to a rather limited yield of target NGs. Conventional surface-assisted synthetic routes often involve a series of cascade reactions with competing reaction pathways, which inevitably lead to the formation of numerous undesired products. The limited yield of target products imposes a formidable challenge toward the practical applications of the NGs synthesized on the surface, because feasibility of the large-scale (i.e., high-yield) fabrication is

vital for their technological relevance. To overcome this issue, it is required to devise a new synthetic route with a new conceptual design of precursor molecules, which allows for their robust transformation toward target NGs via suppressing the side reactions, including undesired intermolecular coupling and intramolecular transformations.

To this end, we design a novel surface-assisted synthetic route for the ultrahigh-yield fabrication of a hexagonal zigzag-edged NG, namely, circumcoronene, on the Cu(111) surface. Our surface-assisted cyclodehydrogenation strategy relies on the robust dehydrogenative coupling of the methyl groups at the adjacent sites of the rationally designed precursor molecule **II**, followed by dehydrogenative aromatization on metallic substrate (Fig. 1A). Because of a remarkably high-yield synthesis and elegant molecule-substrate interaction, circumcoronene molecules self-assemble into extended superlattices (Fig. 1D). The structural and electronic properties of both individual molecules and their superlattices are characterized via bond-resolved noncontact atomic force microscopy (ncAFM) and scanning tunneling microscopy (STM)/scanning tunneling spectroscopy (STS), respectively. We demonstrate that the circumcoronene superlattice confines the two-dimensional (2D) electron gas on Cu(111) into an extended Kagome-honeycomb electronic lattice (Fig. 1D) with a two-flat-band electronic structure, in contrast to a conventional electronic Kagome lattice with a single dispersionless band. Our findings open a new route for both ultrahigh-yield synthesis of elusive NGs and precise fabrication of artificial 2D lattices with possible nontrivial electronic properties toward future technological applications.

RESULTS AND DISCUSSION

Surface-assisted synthesis of circumcoronene

To fabricate a circumcoronene, we adopt a strategy of a surface-assisted cyclodehydrogenation of the rationally designed precursor molecule **II** (C₅₄H₄₂) that was obtained via multistep wet-chemistry synthesis (Fig. 1A). Oxidative cyclodehydrogenation

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Fig. 1. On-surface synthesis of circumcoronene-based superlattice. (A) Precursor molecule **II** undergoes cyclodehydrogenation on Cu(111) via methyl radical-radical coupling toward formation of the circumcoronene. The cyclodehydrogenation results in ultrahigh-yield circumcoronene molecules. (B) STM image acquired after deposition of the precursor **II** onto Cu(111) at room temperature. Inset shows the STM image of an individual precursor molecule. (C) STM image collected after sample annealing at 140°C, revealing the presence of the circumcoronene-based self-assembly on Cu(111). (D) Schematic illustration of the chiral Kagome-honeycomb lattice.

of the hexakis(3,5-dimethylphenyl)benzene precursor molecule **I** by iron(III) chloride gave the hexa-*peri*-hexabenzocoronene (HBC) precursor molecule **II** carrying 12 methyl groups at the bay regions in nearly quantitative yield. This additional step reduces the rotational degree of freedom of six pairs of methyl groups and locks their positions ready for the subsequent dehydrogenative coupling into circumcoronene with zigzag-edged topology on the catalytically active Cu(111) surface. Such a precursor design prevents the undesired intramolecular and intermolecular side reactions, allowing for high-yield synthesis of zigzag-edged NGs. Furthermore, the presence of methyl groups in precursor **II** reduces intermolecular π - π interactions within the compound and thus facilitates in situ sublimation for on-surface synthesis. A complete synthetic route of **II** along with standard structural characterization is presented in the Supplementary Materials (note S1). The conformation and self-assembly of this molecule on Ag(111) and Au(111) were previously investigated (19, 20), but its chemical reaction on metal surface was not reported.

Precursor **II** was then deposited onto the Cu(111) surface held at room temperature via thermal sublimation from a Knudsen cell under ultrahigh vacuum (UHV) conditions. Upon deposition onto the Cu(111) surface, precursor molecules self-assemble into close-packed ordered structures (Fig. 1B). Precursor **II** adopts a nonplanar saddle-shape conformation geometry (20, 21), which results in the triangular-shaped topography in STM images (inset in Fig. 1B). Subsequently, the transformation of precursor **II** into the target circumcoronene was achieved by thermal annealing at ~140°C for 1 hour, whereby the intramolecular cyclodehydrogenation of the methyl groups can be triggered. The chemical transformation of precursor **II** into the final product involves two major steps, namely, the dehydrogenative coupling of the adjacent methyl groups followed by dehydrogenative aromatization on Cu(111) (fig. S5), leading to a

remarkably high product yield (>98%) (Fig. 1C), in stark contrast to the case of precursor **I** (unsuccessful, fig. S7). Note that the product yield is determined by a statistical analysis of multiple large-area STM images (~2000 nm²) (refer to fig. S6 for details).

Characterization of single circumcoronenes and their superlattice

A high-yield synthesis allows individual circumcoronene molecules to self-organize into extended superlattices characterized by a $\sqrt{3} \times \sqrt{3}$ elementary unit cell with a lattice constant of 18.4 Å (Fig. 2A). It is noted that the average-sized superlattice domain contains ~330 molecules, which occupy a total area of ~1000 nm² (the largest domains observed consist of ~520 molecules). To determine an accurate intramolecular structure of individual molecules and their superlattices, we conducted a constant-height ncAFM measurement with a carbon monoxide (CO) terminated tip (22–26). The frequency shift (Δf) contrast revealed by constant-height ncAFM images reflects the magnitude of short-range forces acting between CO-terminated tip and the sample in a Pauli repulsion regime. Therefore, bright areas with a higher Δf value represent regions with an enhanced electron density. A set of ncAFM images collected at various tip-sample heights over individual circumcoronenes (Fig. 2, B and C) unambiguously reveal the molecular skeleton consisting of 19 fused benzene rings with six zigzag edges. Close-up image reveals that the C–C bonds at the molecule's corners exhibit the shortest length when imaged at small tip-sample separations (Fig. 2C). At a moderate tip-sample distance (Fig. 2B), the corner bonds also show a brighter ncAFM contrast, suggesting a higher bond order of these C–C bonds as compared to the other ones. The calculation of the free-standing circumcoronene based on the B3LYP DFT (density functional theory) functional unambiguously reveals a localization of the shortest C–C bonds (or bonds with highest order) at the molecular corners (Fig. 2H). Such

Fig. 2. Bond-resolved ncAFM imaging of the circumcoronene superlattice. (A) Large-scale ncAFM image of the circumcoronene array, collected with a CO-decorated tip. Defective circumcoronene molecule (indicated by yellow arrow) constitutes an approximately 2% of total products. (B and C) Constant-height ncAFM images of individual circumcoronene molecule, collected at tip-sample distances $z = 40$ pm in (B) and $z = 60$ pm in (C) with respect to set point $V = 20$ mV; $I = 1$ nA. (D and E) ncAFM images simulated using a numerical probe-particle approach, for a physical tip-sample distance of 0 pm in (E) and 20 pm in (D) with respect to set point 14.4 Å. (F) DFT-relaxed model of the circumcoronene-based superlattice on Cu(111). (G) Top view and side view of the single circumcoronene molecule on Cu(111). (H) DFT (B3LYP) calculations of the bond lengths within the circumcoronene molecule. Scale bars (B to E), 500 pm.

a peculiar bond localization within circumcoronene can be described on the basis of Clar's aromatic sextet rule that predicts the formation of the maximal number of aromatic sextets (seven in circumcoronene) with a C=C localization at the molecular corners. ncAFM contrast of peripheral C—C bonds in circumcoronene associated with an increased bond localization of aromatic rings also resembles that of a Kekulene molecule (27).

To uncover the atomic registry of the circumcoronene-based superlattice on Cu(111), we imaged both individual molecules within the superlattice and adjacent Cu(111) lattice, respectively (fig. S8). Superimposing the Cu(111) lattice over the circumcoronene array reveals the adsorption sites of all the circumcoronene molecules corroborated by the DFT-predicted model (Fig. 2F). Furthermore, DFT calculations also unveil that circumcoronene adopts a slightly bent geometry, with peripheral benzene rings pulled toward the underlying Cu(111) surface (Fig. 2G). Such a bent structural conformation of circumcoronene is manifested by a brighter contrast of the central benzene ring in the constant-height ncAFM images. The simulated ncAFM images of DFT-relaxed structures using a modified probe-particle model (28, 29) reproduce all the key features of experimental ncAFM images including a tip-sample-dependent Δf contrast over the corner C—C bonds with the highest bond order and higher Δf contrast of the central benzene ring attributed to the nonplanar geometry of the circumcoronene (Fig. 2, D and E).

Electronic properties of the circumcoronene-based superlattice

We then probed the electronic properties of individual circumcoronene molecules and their superlattice via performing differential

conductance spectroscopy (dI/dV) measurements (Fig. 3 and fig. S9). In contrast to the rather featureless dI/dV spectrum acquired over the bare Cu(111) surface (only Shockley surface state induced peak around -0.4 V), the dI/dV curve collected over circumcoronene at the periphery of the superlattice (blue curve in Fig. 3A) exhibits three characteristic peaks, located at energies of -1.6 V, $+0.9$ V, and $+1.8$ V, respectively, resembling that acquired over an isolated molecule (green curve in Fig. 3A). The dI/dV curve acquired over the circumcoronene in the interior of the superlattice (red curve in Fig. 3A) reveals additional prominent peaks centered at $+0.36$ V (labeled as FB_1) and $+1.15$ V (labeled as FB_2). We further probed the spatial distribution of these electronic states via dI/dV mapping with a CO-functionalized tip. The dI/dV maps collected at the corresponding energies of -1.6 V, $+0.9$ V, and $+1.8$ V reveal the characteristic nodal patterns localized over the molecular backbone, which can be assigned to the highest occupied molecular orbital (HOMO), lowest unoccupied molecular orbital (LUMO), and second lowest unoccupied molecular orbital (LUMO+1), respectively, as supported by their excellent resemblance with the calculated dI/dV maps of these frontier orbitals using a CO-decorated tip (30, 31) (insets in Fig. 3, C, E, and G). In contrast, dI/dV maps collected at energies of the additional electronic states (FB_1 and FB_2) within the circumcoronene superlattice reveal remarkable intensity in between molecules (Fig. 3, D and F), which cannot be explicitly ascribed to tunneling into any frontier molecular orbitals. Furthermore, FB_1 and FB_2 states largely attenuate at the edge of the circumcoronene superlattice as evident from the 2D contour plot of dI/dV spectra acquired across the superlattice (Fig. 3B).

Flat bands in the chiral electronic Kagome-honeycomb lattice

We performed DFT calculation combined with tight-binding (TB) model calculations to understand the physical origin of both FB_1 and FB_2 states in the circumcoronene superlattice. These new electronic states arise from the formation of the circumcoronene superlattice that confine 2D electron gas (2DEG) of the Cu(111) surface into the chiral electronic Kagome-honeycomb lattice. The calculated electrostatic potential landscape over the circumcoronene superlattice shows a notable accumulation of the negative charge in the plane of circumcoronene (Fig. 4B). This negative electrostatic potential is balanced by a positive potential accumulated in the Cu(111) surface regions between molecules. In addition, a comparison of the calculated DOS (density of states) plots of the free-standing circumcoronene and Cu(111)-supported circumcoronene shows that the latter exhibits a downshift of the energetic positions of the molecular orbitals (figs. S10 and S11), attributed to the electron transfer from the Cu(111) substrate to the circumcoronene. Each molecule hosts a negative electrostatic potential and thereby acts as a hexagonally shaped repulsive barrier for creating a periodic modulation and confinement of the 2DEG (32–34). Notably, our DFT calculations also reveal a non-negligible out-of-plane buckling of Cu atoms in the topmost layer of the Cu(111) regions forming a periodic buckling pattern (Fig. 4C) reminiscent of the chiral Kagome-honeycomb lattice (Fig. 4A). Because of a computational feasibility, the semi-infinite Cu(111) surface is represented by only three Cu layers in our DFT calculations, which therefore cannot provide an accurate description of the electronic structures of surface state and the confined states observed experimentally.

The unique electronic structures of the artificial electronic lattice can be well described using the TB model that involves the electron

Fig. 3. Electronic structure of the circumcoronene array. (A) dI/dV spectra collected over the bare Cu(111) substrate (gray), a circumcoronene molecule residing at the periphery of the superlattice (blue), a circumcoronene molecule located in the interior of the superlattice (red), and an isolated molecule (green). a.u., arbitrary units. (B) 2D contour plot of dI/dV spectra acquired along the line indicated in the inset in (A). (C to G) Topography STM image (1 V; 1 nA) and corresponding dI/dV maps collected at energetic positions of states HOMO (C), FB₁ (D), LUMO (E), FB₂ (F), and LUMO+1 (G). Inset images in (C), (E), and (G) are the calculated dI/dV maps of the HOMO (C), LUMO (E), and LUMO+1 (G) of the circumcoronene using a CO-decorated tip.

Fig. 4. Origin and electronic structure of the chiral Kagome-honeycomb lattice. (A) Schematic illustration of the chiral Kagome-honeycomb lattice. (B) Electrostatic potential of the circumcoronene superlattice. (C) The out-of-plane buckling of Cu atoms in the topmost Cu(111) layer. (D and E) A k-space dispersion of the electronic bands of the Kagome and chiral Kagome-honeycomb lattices, calculated using the TB model for hopping constants $t_1 = t_2 = t_3 = -0.5$ eV (D) and $t_1 = t_2 = t_3 = t_4 = -0.5$ eV (E). (F) Experimental dI/dV spectra reveal the existence of two characteristic peaks (FB₁ and FB₂) associated with two flat bands of the chiral electronic Kagome-honeycomb lattice.

hopping strength between atomic orbitals associated to lattice sites. TB band structures of a standard Kagome and Kagome-honeycomb lattice are shown in Fig. 4 (D and E, respectively). The common Kagome lattice is composed of a network of corner-sharing triangles, leading to the destructive phase interference of electron hopping pathways between nearest-neighbor lattice sites. As such, it precludes electron hopping between neighboring unit cells and results in the localization of electronic state within a single Kagome hexagon (35, 36). Such a unique real-space localization of the electronic state further manifests itself by the formation of the empty state dispersionless band (i.e., flat bands) alongside Dirac bands in a reciprocal space, as revealed by TB band structure derived from hopping parameters of $t_1 = t_2 = t_3 = -0.5$ eV. In contrast, the Kagome-honeycomb lattice can be virtually derived from the Kagome lattice by introducing an additional bond between frustrated Kagome's triangles, associated with the additional hopping constant t_4 . The TB band structure of the Kagome-honeycomb lattice (Fig. 4E) with hopping constants of $t_1 = t_2 = t_3 = t_4 = -0.5$ eV exhibits two flat bands below and above Dirac bands, which can be attributed to a pair of the intense peaks (FB_1 and FB_2) in the electronic spectrum of the circumcoronene-based superlattice on Cu(111) (Fig. 4F).

The electronic Kagome-honeycomb lattice emerges as “antilattice” to the circumcoronene superlattice, whereby individual molecules act as a hexagonal repulsive barrier for 2DEG. The underlying physical picture is remarkably similar to the artificial electronic lattices (37–43) patterned by CO molecules on the Cu(111) surface via STM manipulation, wherein each CO molecule also acts as a repulsive scatter for 2DEG. Although STM manipulation endows an ultimate control over the lattice topology, this approach is unsuitable for practical applications due to a lack of scalability. Our findings surpass this limitation, paving a viable route toward a scalable bottom-up fabrication of the extended 2D electronic lattices, where NG with tailored topology is used as a “building block.” We envisage that a rational design of NGs with different size and edge topologies can potentially serve as a versatile “tuning knob” for the efficient modulation of the hopping strength and electronic structure of artificially engineered electronic lattices.

We have demonstrated a conceptually new on-surface synthesis route toward the ultrahigh-yield fabrication of the hexagonal-shaped circumcoronene molecule on Cu(111) via surface-assisted intramolecular cyclodehydrogenation of the rationally designed precursor molecule. A simultaneous assembly of large-area circumcoronene array results in a periodic confinement of the 2DEG into a unique chiral electronic Kagome-honeycomb lattice associated with the peculiar electrostatic potential landscape of the circumcoronene array on Cu(111). Our results open up a new avenue for high-yield fabrication of atomically precise elusive NGs with desired topologies toward large-scale fabrication of artificial 2D lattices with possible nontrivial electronic properties. The 2D Kagome-honeycomb lattices with a flat-band electronic structure can be particularly interesting in a wide scope of condensed matter physics, because of their prosperous potential in the realization of a variety of exotic many-body phenomena, including anomalous quantum Hall states (44–46), Wigner crystallization (47), and topological insulating transitions (48).

MATERIALS AND METHODS

The qPlus STM/ncAFM measurements

The STM and ncAFM experiments were performed under UHV conditions at 4.4 K using a commercial Omicron LT STM/AFM

machine. We used the commercial qPlus sensor (49) with a resonant frequency of $f_0 = 28$ kHz, a stiffness of $k = 1800$ N/m, and a quality factor of 11,500, which was operated in the frequency-modulation mode with an oscillation amplitude of $A = 100$ pm. All simultaneous STM/ncAFM images were collected in constant height mode. The tip apex was routinely decorated by a CO molecule by its controllable picking from the bare Cu(111) surface.

Sample preparation

Cu(111) single crystal (MaTeck GmbH) was cleaned by multiple cycles of Ar⁺ sputtering and annealing. The precursor molecules were deposited from the Knudsen cell (MBE-Komponenten GmbH) at 350°C onto the clean Cu(111) surface under UHV conditions (base pressure $< 3 \times 10^{-9}$ mbar). After the deposition of the precursor molecules, the sample was annealed at 140°C for 1 hour for the fabrication of circumcoronene arrays. Subsequently, the sample was transferred into the STM/AFM head held at a temperature of 4.5 K for ncAFM/STM imaging and characterization.

DFT simulations

DFT calculations were performed using the FHI-aims code (50) and Perdew-Burke-Ernzerhof generalized gradient approximation of the exchange-correlation potential (51, 52) to obtain the electronic properties of the circumcoronene molecule on the Cu(111) surface. The Tkatchenko-Scheffler treatment of the van der Waals (vdW) interactions was used during the calculations (53). We modeled the Cu(111) surface by a 3×3 supercell composed of three layers. After placing the molecule on the surface, we performed structural relaxations of all the atoms of the slab, except the last Cu layer. The optimization procedure was stopped when the total energy and the remaining atomic forces were found below 10^{-5} eV and 10^{-2} eV/Å, respectively. A Monkhorst-Pack grid of $3 \times 3 \times 1$ was used to integrate the Brillouin zone. Furthermore, the calculated total charge density and the Hartree potential were used to analyze the electronic interactions between the circumcoronene and the surface.

TB calculations

We carried out calculations of a TB model using the nearest-neighbor interaction between nodes of the Kagome-honeycomb lattice shown in Fig. 4E. All on-site energies have been set to zero and the nearest-neighbor hopping was set to $t = -0.5$ eV to mimic the experimental evidence. To obtain the band plot shown in Fig. 4E, TB Hamiltonian was diagonalized along a dense set of high-symmetry k-points.

SPM simulations

AFM image simulations were performed on the basis of the Probe Particle code (28, 29), which includes vdW and electrostatic interactions between the CO tip and the surface. Lennard-Jones potential was used for the treatment of Pauli repulsion and vdW interactions. The Hartree potential obtained from DFT calculations was used to calculate the electrostatic forces. All the simulations were done with an effective atomic radius $R_c = 1.661$ Å and a lateral stiffness of $k = 0.25$ N/m. dI/dV maps were simulated using the PP-STM code (30) with a CO rigid tip (without tip-apex relaxation). We first calculated the eigenstates of the circumcoronene orbitals (with the FHI-aims code) that were subsequently used in the simulations. The dI/dV maps were obtained at characteristic energy positions corresponding to frontier orbitals (−1.6 V, +0.9 V, and +1.8 V).

Note that the Lorentzian width parameter for broadening of the eigenstates was set to be 0.001 eV.

SUPPLEMENTARY MATERIALS

Supplementary material for this article is available at <http://advances.sciencemag.org/cgi/content/full/7/3/eabf0269/DC1>

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