Published in final edited form as:

Phys Rev E Stat Nonlin Soft Matter Phys. 2013 December ; 88(6): 062818. doi:10.1103/PhysRevE.

88.062818

Multi-opinion coevolving voter model with infinitely many phase transitions

Feng Shi¹, Peter J. Mucha¹, and Richard Durrett²

¹Dept. of Mathematics, CB#3250, U. of North Carolina, Chapel Hill, NC 27599-3250

²Dept. of Mathematics, Box 90320, Duke University, Durham, NC 27708-0320

Abstract

We consider an idealized model in which individuals' changing opinions and their social network coevolve, with disagreements between neighbors in the network resolved either through one imitating the opinion of the other or by reassignment of the discordant edge. Specifically, an interaction between x and one of its neighbors y leads to x imitating y with probability $(1-\alpha)$ and otherwise (i.e., with probability α) x cutting its tie to y in order to instead connect to a randomly chosen individual. Building on previous work about the two-opinion case, we study the multiple-opinion situation, finding that the model has infinitely many phase transitions (in the large graph limit with infinitely many initial opinions). Moreover, the formulas describing the end states of these processes are remarkably simple when expressed as a function of $\beta = \alpha/(1-\alpha)$.

I. INTRODUCTION

In the last decade, there have been a number of studies of systems in which the states of individuals and the connections between them coevolve, see [1, 2]. The systems considered include evolutionary games [3]–[7] and epidemics [8]–[12], but here we will concentrate on the spread of opinions [13]–[18]. A natural question to ask is that whether individuals develop opinions based on the opinions of their neighbors, or whether they form connections because they share similar opinions. Holme and Newman [19] proposed a simple physical model (referred to as the ‘coevolving voter model’) combining these two processes. Different from the models of cascades [20]–[22] which are also widely used in the study of opinion spread, the coevolving voter model allows an agent to switch between different opinions and the network topology to change accordingly, yet we assume that agents impose equal influence over each other in contrast to multi-state complex contagions [23]–[25]. This model provides building blocks to quantitatively study collective behaviors in various social systems, e.g., segregation of a population into two or more communities with different political opinions, religious beliefs, cultural traits, etc.

Our goal in the present work is to further explore the remarkably complex behavior present in linear voter model systems, generalizing the model proposed by Holme and Newman [19]. In their model there is a network of N vertices and M edges. The individual at vertex v has an opinion $\xi(v)$ from a set of G possible opinions and the number of people per opinion $\gamma_N = N/G$ stays bounded as N gets large. On each step of the process, a vertex x is picked at random. If its degree $d(x)$ equals 0, nothing happens. If $d(x) > 0$, (i) then with probability $1 -$

α a random neighbor y of x is selected and we set $\xi(x) = \xi(y)$; (ii) otherwise (i.e., with probability α) an edge attached to vertex x is selected and the other end of that edge is moved to a vertex chosen at random from those with opinion $\xi(x)$. This process continues until the ‘freezing time’ τ , at which there are no longer any discordant edges—that is, there are no edges connecting individuals with different opinions.

For $\alpha = 1$, only rewiring steps occur, so once all of the M edges have been touched, the graph has been disconnected into G components, each consisting of individuals who share the same opinion. Since none of the opinions have changed, the components of the final graph are all small (i.e., their sizes are Poisson with mean γ_N). By classical results for the coupon collector’s problem, this requires $\sim M \log M$ updates, see e.g., page 57 in [26]. In the case of sparse graphs we consider here $M \sim cN$ (i.e., $M/N \rightarrow c$) so the number of steps is $O(N \log N)$, i.e., when N is large it will be $\approx CN \log N$.

In contrast, for $\alpha = 0$ this system reduces to the voter model on a static graph. If we suppose that the initial graph is an Erdős-Rényi random graph with average degree $\lambda > 1$, then (see e.g., Chapter 2 of [27]) there is a “giant component” that contains a positive fraction, μ , of the vertices and the second largest component is small having only $O(\log N)$ vertices. The voter model on the giant component will reach consensus in $O(N^2)$ steps (see, e.g., Section 6.9 of [27]), so the end result is that one opinion has μN followers while all of the other groups are small.

Using simulation and finite size scaling, Holme and Newman showed that there is a critical value α_c so that for $\alpha > \alpha_c$ all of the opinions have a small number of followers at the end of the process, while for $\alpha < \alpha_c$ “a giant community of like-minded individuals forms.” When the average degree $\lambda = 2M/N = 4$ and the number of individuals per opinion $\gamma_N \rightarrow 10$, this transition occurs at $\alpha_c \approx 0.46$. See [28]–[31] for recent work on this model.

In [32], we studied a two-opinion version of this model in which on each step an edge is chosen at random and is given a random orientation, (x, y) . If the individuals at the two ends have the same opinion nothing happens. If they differ, then (i) with probability $1 - \alpha$ we set $\xi(x) = \xi(y)$; (ii) otherwise (i.e., with probability α) x breaks its edge to y and reconnects to (a) a vertex chosen at random from those with opinion $\xi(x)$, a process we label ‘rewire-to-same’, or (b) at random from the graph, a process we label ‘rewire-to-random’. Here, we will concentrate on the second rewiring option, rewire-to-random. While this process may be less intuitive than the rewire-to-same version, it has a more interesting phase transition, as documented in [32].

The remainder of this paper is organized as follows. In Section II, we recall the main results from [32] that provide context for our observations of the multiple-opinion case, which we begin to explore in Section III. We then continue in Section IV with further quantitative details about the phase transitions and their underlying quasi-stationary distributions, before concluding comments in Section V.

II. TWO-OPINION MODEL

Suppose, for concreteness, that the initial network is an Erdős-Rényi random graph in which each individual has average degree $\lambda > 1$, and that vertices are assigned opinions 1 and 0 independently with probabilities u and $1 - u$. For concreteness, all of the numerical results presented in the rest of the main body of the paper are for $\lambda = 4$, with additional results presented in the Supplemental Materials [34] for $\lambda = 2$ and 8 confirming the qualitative phenomena described here. We in particular note that the initial average degree does not change the qualitative behavior of our model, yet the quantitative details such as the critical value α_c depend on the value of λ . Additionally, our results are not limited to Erdős-Rényi random graphs. As shown below, the phase transition is due to the change in the balance between the processes of opinion change and rewiring, which together dictate the properties of the resulting networks, which deviate from Erdős-Rényi as the dynamics proceed. Similarly, other sparse initial graph conditions (subject to the average degree, λ , which is a conserved quantity here) will converge to the same ensemble of random graphs under this dynamics. Initial network structures affect the details of those dynamics but the qualitative behavior of the model, determined by the competition of the two processes, will be the same. Simulations suggest that the system has the following

Phase transition

For each initial density $u \in [0, 1/2]$ there is a critical value $\alpha_c(u)$ so that for $\alpha > \alpha_c(u)$, freezing occurs after $O(N \log N)$ updates and the fraction of voters with the minority opinion at the end is $\rho(\alpha, u) \approx u$. For $\alpha < \alpha_c(u)$ freezing is slow, requiring $O(N^2)$ updates, and $\rho(\alpha, u) \approx \rho(\alpha, 0.5)$.

Different from the phase transition studied in [19], the phase transition introduced above is not a classical physical phase transition from disorder to order, e.g., from a state with diverse opinions to a consensus state, but a point where there is a change in the qualitative behavior of the model. Below criticality, the system quickly converges to a quasi-stationary distribution and diffuses until freezing occurs after $O(N^2)$ updates, while above criticality freezing occurs rapidly after $O(N \log N)$ updates and the densities of opinions at the end is the same as their initial densities. The phase transition is due to the change in the balance between the processes of opinion change and rewiring (please see the Supplemental Materials [34] for details). To help understand the last statement in the definition of the phase transition, the reader should consult the picture in Figure 1. If the initial fraction of 1's $u = 1/2$ then as α decreases from 1, the ending density $\rho(\alpha, 1/2)$ stays constant at 1/2 until $\alpha = \alpha_c(1/2)$ and then decreases to a value close to 0 at $\alpha = 0$. For convenience, we call the graph of $\rho(\alpha, 1/2)$ for $\alpha < \alpha_c \doteq 0.74$, the *universal curve*. If the initial density is $u < 1/2$, then the ending density $\rho(\alpha, u)$ stays constant at u until the flat line (α, u) hits the universal curve and then $\rho(\alpha, u) \approx \rho(\alpha, 0.5)$ for $\alpha < \alpha_c(u)$. Figure 2 shows the average freezing time (i.e., number of updates needed to reach the frozen configuration) at each α . Below the critical α freezing takes almost $O(N^2)$ updates while it only requires $O(N \log N)$ updates above criticality, and where the transition occurs agrees with critical α 's in Figure 1. The main aim of [32] was to use simulations, heuristic arguments, and approximate models to explain the presence and properties of this universal curve describing the freezing states that

result from the slow-freezing process. To make it easier to compare the results here with the previous paper, we rescale time so that times between updating steps are exponential with rate M , where M is the total number of edges. Let $\xi_i(x)$ be the opinion of the voter at x at time t . Analogues to the voter model on a static finite torus, for the two-opinion coevolving voter model we find the following (see [32] for details)

Quasi-stationary distributions

Let $v(a) = \rho(a, 0.5)$. If $a < a_c(1/2)$ and $v(a) < u - 1/2$ then starting from product measure with density u of 1's (i.e., $\xi_0(x)$ are independent and equal to 1 with probability u), $\xi_t(x)$ converges rapidly to a quasi-stationary distribution $\nu_{a,u}$. At time tM the coevolving voter model looks locally like $\nu_{a,\theta(t)}$ where the density $\theta(t)$ changes according to a generalized Wright-Fisher diffusion process

$$d\theta_t = \sqrt{(1-\alpha)[c_\alpha\theta_t(1-\theta_t) - b_\alpha]}dB_t \quad (1)$$

until θ_t reaches the two absorbing states $v(a)$ or $1-v(a)$, the two solutions of $c_\alpha x(1-x) = b_\alpha$.

In (1) B_t is a standard Brownian motion and c_α and b_α are unknown coefficients of the ‘arch’ studied in the following sections. To begin to explain the behavior of θ_t given in (1), note that when an edge is picked with two endpoints that differ, a rewiring will not change the number of 1's, while a voting event, which occurs with probability $(1 - \alpha)$, will result in an increase or decrease of the number of 1's with equal probability. Thus the quantity under the square root is $(1-\alpha)N_{01}/M$. When $\theta_t = u$ the rate at which 0–1 edges are chosen is equal to the expected fraction of 0–1 edges under $\nu_{a,u}$, which is $c_\alpha u(1 - u) - b_\alpha$.

To further explain the phrase “quasi-stationary distributions” in this context, we refer the reader to Figure 3. Let $N_1(t)$ be the number of vertices in state 1 at time t , $N_{01}(t)$ be the number of 0–1 edges (that is, the number of edges connecting nodes x and y with $\xi(x) = 0$, $\xi(y) = 1$). Similarly, let $N_{abc}(t)$ be the number of connected triples x - y - z with $\xi(x) = a$, $\xi(y) = b$, and $\xi(z) = c$. The top panel of Figure 3 plots $N_{01}(t)/M$ versus $N_1(t)/N$ for five different simulations (with different initial densities, u) for $\alpha = 0.5$. Note that in each case the simulation rapidly approaches a curve $\approx 1.710x(1 - x) - 0.188$ and then diffuses along the curve until freezing is reached ($N_{01} = 0$). At both of the possible freezing points on the curve, the fraction of the minority opinion is ≈ 0.12 , in accordance with the simulation in Figure 1.

The bottom panel of Figure 3 similarly plots $N_{010}(t)/N$ versus $N_1(t)/N$ for $\alpha = 0.5$ and $u = 1/2$. Again the simulation rapidly approaches a curve (approximately cubic) and diffuses along it until freezing is reached. Since $N_{010} = 0$ if $N_{01} = 0$, and it is very unlikely that all 0–1's only occur in 0-1-1 triples, the zeros of the cubic curve for 0-1-0 and quadratic curve for 0–1 coincide.

One can repeat the simulations in Figure 3 for other network measurements, with the result that their values are similarly determined by the density $u(t) = N_1(t)/N$. This is somewhat

analogous to a stationary distribution from equilibrium statistical mechanics—e.g., the Maxwell-Boltzmann distribution associating the velocity distribution with the temperature. We call our distributions quasi-stationary because our system is a finite state Markov chain, which will eventually reach one of its many absorbing states $N_{01} = 0$, and hence there is no true stationary distribution. Nevertheless, an improved understanding of the system is obtained from these observations, displaying a fast dynamics rapidly converging to a family of neutrally-stable quasi-stationary distributions followed by slow, diffusive dynamics through the space local to the quasi-stationary distributions until freezing is reached.

As shown in [32], the behaviors for the rewire-to-same model in terms of quasi-stationary distributions are very similar, but with small differences from the rewire-to-random model that yield fundamentally different freezing states. In rewire-to-same, there are quasi-stationary distributions $v'_{\alpha,u}$ under which the expected fraction of 0–1 edges is $c'_\alpha u(1-u)$. Again the simulation comes rapidly to this curve and diffuses along it until freezing is reached. That is, unlike Figure 3 (Top), the arches of quasi-stationary N_{01}/M values versus N_1/N maintain their zeros at $N_1/N = \{0, 1\}$. Thus, for $\alpha < \alpha'_c(1/2)$, the minority fraction obtained at the freezing time is always ≈ 0 for rewire-to-same.

III. MULTI-OPINION MODELS

Böhme and Gross [33] have studied the three-opinion version of the coevolving voter model with rewire-to-same dynamics. In this case, the limiting behavior is complicated – one may have partial fragmentation (1's split off rapidly from the 2's and 3's) in addition to full fragmentation and coexistence of the three opinions. See their Figures 3–5. As we will see in the present section, the behavior of the multi-opinion rewire-to-random model is much simpler because small groups of individuals with the same opinion will be drawn back into the giant component. We thus aim to extend the understanding of the two-opinion model behavior to larger numbers of opinions.

Consider now the k -opinion model in which voters are assigned independent initial opinions that are equal to i with probability u_i . Let $u = (u_1, u_2, \dots, u_k)$ and let N be the number of edges at which the endpoint opinions differ. When $k = 3$, frequencies of the three types must lie in the triangle of possible values $\mathcal{U} = \{u = (u_1, u_2, u_3) : u_i \geq 0, \sum_i u_i = 1\}$. To preserve symmetry, we draw $\mathcal{U}$ as an equilateral triangle in barycentric coordinates by mapping $(x, y, z) \rightarrow (x, z/\sqrt{3}/2)$. The top panel in Figure 4 plots $N(t)/M$ as a function of the opinion densities as the system evolves, generalizing the one-dimensional arch observed for $k = 2$ to a two-dimensional cap for $k = 3$.

Generalizing the parabolic form of the arch for $k = 2$, we conjecture

$$E_u N_{\neq}/M = \frac{c_2(\alpha)}{2} \left(1 - \sum_{i=1}^k u_i^2 \right) - c_0(\alpha), \quad (2)$$

where the expectation E_u of the fraction of discordant edges N/M is taken under the quasi-stationary distribution $\nu_{a,u}$, and $c_2(a)$ and $c_0(a)$ are unknown coefficients to be fitted to the quasi-stationary distributions. As in the two opinion case, the simulated values come quickly to the surface and then diffuse along it. In some situations, one opinion is lost before freezing occurs and the evolution reduces to that for the two opinion case. However, in one of the simulations shown, the realization ending with $x \approx 0.5$, all three opinions persist until the end.

The picture is somewhat easier to understand if we look at the cap from a top view, where the $E_u N = 0$ level sets for different a are observed to be circles. In the bottom panel of Figure 4 we plot the $E_u N = 0$ circles for different a 's fitted from simulation data using (2) as well as the freezing opinion frequencies from the simulations (indicated by small circle data points). The two agree with each other up to small stochastic fluctuations. The size of the $E_u N = 0$ level set then dictates different freezing state properties. For example, the circle corresponding to $a = 0.5$ intersects in three disconnected arcs. As a increases, the radius of the $E_u N = 0$ level set decreases. When $a > a_c(1/2)$, the critical value of the two opinion model, the circle $E_u N = 0$ falls fully inside the triangle, so an initial condition including all three opinions will continue to demonstrate all three opinions at freezing. For example, the small circles around the innermost circle give the ending frequencies for several simulations for $a = 0.8$. If the initial frequencies fall within the $E_u N = 0$ circle, then the model will quickly relax to the quasi-stationary distributions above the circle and then diffuse along the cap until freezing is reached at some $E_u N = 0$ point. If instead the initial frequencies u fall outside the $E_u N = 0$ circle—that is, for a above the phase transition point $a_3(u)$ —the freezing time jumps from $O(N^2)$ to $O(N \log N)$, similar to $a_c(u)$ for the two-opinion model, with the final opinion frequencies essentially the same as the initial u . What is new in this case is that when starting with three opinions and $a_c(u) < a < a_3(u)$, the system always ends up with three distinct opinions.

For $k > 3$, our simulation results indicate the same type of behavior as the system evolves. We define a_k to be the largest a for which freezing takes $O(N^2)$ updates when we start with k opinions with density $1/k$ for each opinion. Then as $k \rightarrow \infty$ the multi-opinion model has infinitely many phase transitions. When $a_k < a < a_{k+1}$, freezing occurs after $O(N \log N)$ steps if we start with k opinions, while if we start with $k+1$ equally likely opinions the system quickly converges to a quasi-stationary distribution and diffuses until freezing occurs after $O(N^2)$ updates and there will always be $k+1$ opinions present at the end. The associated picture is the natural dimensional extension of the relationship between the $k=2$ and $k=3$ models: just as $a_2 = a_c(1/2)$ corresponds to the point at which the $E_u N = 0$ circle for $k=3$ is the inscribed circle within the triangle, a_3 corresponds to the point at which the $E_u N = 0$ circle reaches zero radius—that is, the point at which the $E_u N = 0$ sphere for $k=4$ has become the inscribed sphere within the corresponding barycentric tetrahedron.

IV. QUANTITATIVE CHARACTERIZATION OF QUASI-STATIONARY DISTRIBUTIONS

For each k we simulate our multi-opinion rewire-to-random model starting from k opinions with each opinion taking $1/k$ fraction of nodes at random for a wide range of α 's. Generalizing the picture of the one-dimensional arch for $k=2$ and the two-dimensional cap for $k=3$, the number of discordant edges as a function of frequencies conjectured in (2) is a co-dimension 1 hypersurface characterizing the quasi-stationary states, and the behavior of the equal-initial-populations case will allow us to describe this surface, thereby characterizing behaviors for general initial populations.

First the critical α_k 's are identified when the slow diffusion of N cannot be observed for the first time as α increases from 0 to 1. Then we fit $N(t)/M$ to $u_i(t) = N_i(t)/N$ ($i=1, \dots, k$) using (2) at every α up to α_k , and plot the fitted coefficients c_0 and c_2 against $\beta = \alpha(1-\alpha)$ in Figure 5. Remarkably, the coefficients in (2) appear to be well approximated by linear functions of $\beta = \alpha(1-\alpha)$. The graphs shows some curvature near $\beta=0$ ($\alpha=0$) which may be caused by the fact that $\beta=0$ ($\alpha=0$) corresponds to a voter model without evolution of the underlying network. In the rest of the paper, we will work with β for simplicity. Naturally, critical points α_k translate to $\beta_k = \alpha_k(1-\alpha_k)$.

The fitted coefficients from the 2-opinion model deviate slightly from those fitted from higher-order models, which implies that (2) is not universal for the multi-opinion model and higher-order terms are possible. However, while the discrepancy between the fitted coefficients of the 2-opinion model and those of the 3-opinion one is apparent, difference between fitted coefficients of higher-order models is negligible, which implies that the inclusion of higher-order terms beyond the 3rd would not make significant changes to the equation. To probe the effect of higher-order terms we introduce terms up to k th order for k opinions. Noting $(\sum_j u_j)^2 = 1$, (2) is equivalent to:

$$E_u N_{\neq}/M = -c_0(\alpha) + c_2(\alpha) \sum_{i,j=1; i>j}^k u_i u_j. \quad (3)$$

Given the symmetry of the system in u_i 's, the only possible choice in degree- k polynomials is:

$$\begin{aligned} E_u N_{\neq}/M = & -c_0(\alpha) + c_2(\alpha) \sum_{\{i_1, i_2\} \in \mathcal{A}_2} u_{i_1} u_{i_2} \\ & + c_3(\alpha) \sum_{\{i_1, i_2, i_3\} \in \mathcal{A}_3} u_{i_1} u_{i_2} u_{i_3} + \dots \\ & + c_k(\alpha) \sum_{\{i_1, \dots, i_k\} \in \mathcal{A}_k} u_{i_1} u_{i_2} \dots u_{i_k}, \end{aligned} \quad (4)$$

where $\mathcal{A}_i$ is the collection of all i -element subsets of $\{1, 2, \dots, k\}$. Using the same simulation data as above, we refit $N(t)/M$ to $u_i(t)$'s ($i=1, \dots, k$) according to the generalized formula (4) and plot the fitted coefficients c_0 and c_2 against β in Figure 6. Fitting

diagnostics suggest that higher-order terms are significant from zero (with p -value $< 10^{-4}$) and it can be seen that those terms explain the inconsistency between fitted coefficients of different models in Figures 5. However, the difference between the two fitted functions of (2) and (4) is actually small ($\approx .1$ in L^2 -norm) and thus higher-order terms are small corrections to the hyper-surface (2).

Values of the coefficients $c_i(\beta)$ for the three opinion model near its critical value $\beta_3 \approx 5.2$ show some scatter, but this is to be expected since the surface is very small at this point. Values for the four opinion model appear to become more difficult to fit prior to β_4 since $E_u N = 0$ is a three-dimensional hyper-surface in four-dimensional space, so much more data is required to get reliable estimates of coefficients.

As is visually apparent in Figure 6, the coefficients c_0 and c_2 for the first two terms in (4) are well approximated by linear functions, with best fits $c_0(\beta) \approx 0.22\beta$ and $c_2(\beta) \approx 1.40 + 0.36\beta$. In contrast, coefficients for higher-order terms are not linear in β (e.g., see Figure 7 for $c_3(\beta)$). Moreover, the fitted coefficients for higher-order terms become particularly noisy as β increases (again, see Figure 7) at the same time that their relative contribution to the shape of the hyper-surface decreases (relative to the linearly increasing c_0 and c_2). For comparison, the best fits for c_0 and c_2 in (2) (as in Figure 5) are

$$c_2(\beta) \approx 1.35 + 0.52\beta, \quad c_0(\beta) \approx 0.26\beta. \quad (5)$$

Since (2) well approximates the higher-order hyper-surface (4), its simple form can be used to estimate the critical points for phase transitions. Combining (2) and (5), and then solving

$$\frac{1}{2}(1.35 + 0.52\beta)(1 - k(1/k)^2) - 0.26\beta = 0$$

gives

$$\beta_k \approx 2.6(k-1).$$

which agrees with the critical β_k 's identified when the slow diffusion of N cannot be observed in simulations as β increases.

V. CONCLUSION

Our multi-opinion voter model has infinitely many phase transitions as $k \rightarrow \infty$. When $\beta_k < \beta < \beta_{k+1}$, freezing occurs rapidly when we start with k opinions; however, starting with $k + 1$ equally likely opinions will always yield $k + 1$ opinions present at freezing for this β range (Figure 8). While all of the numerical results presented above are for $\lambda = 4$, additional results presented in the Supplemental Materials [34] for $\lambda = 2$ and 8 confirm the generality of the qualitative phenomena. To a good approximation, $\beta_k \approx 2.6(k - 1)$ (for average degree

$\lambda = 4$), but the departures from linearity in the plots of $c_2(\beta)$ and $c_0(\beta)$ suggest that this result is not exact. However, formulas for various quantities associated with this model are close to polynomials, so an exact solution may be possible.

Acknowledgments

RD was supported by the National Science Foundation (DMS-1005470). FS and PJM were supported by the National Science Foundation (DMS-0645369) and the National Institute of General Medical Sciences (Award Number R21GM099493). The content is solely the responsibility of the authors and does not necessarily represent the official views of the agencies supporting this work.

References

1. Gross T, Blasius B. Adaptive coevolutionary networks: a review. *J Royal Soc Interface*. 2008; 5:259–271.
2. Gross, T.; Sayama, H. *Adaptive Networks*. Springer; Heidelberg: 2009.
3. Ebel H, Bornholdt S. Coevolutionary games on networks. *Phys Rev E*. 2002; 66:056118.
4. Zimmerman MG, Eguíluz VM, San Miguel M. Coevolution of dynamical states and interactions on networks. *Phys Rev E*. 2004; 69:065102.
5. Zimmerman MG, Eguíluz VM. Cooperation, social networks, and the emergence of leadership in a prisoner's dilemma game. *Phys Rev E*. 2005; 72:056118.
6. Santos FC, Pacheco JM, Lenaerts T. Cooperation prevails when individuals adjust their social ties. *PLoS Comp Biol*. 2006; 2(10): e140.
7. Pacheco JM, Traulsen A, Nowak MA. Coevolution of strategy and structure in complex networks with dynamical linking. *Phys Rev Letters*. 2006; 97:258103.
8. Gross T, Dommar D'Lima CJ, Blasius B. Epidemic dynamics on an adaptive network. *Phys Rev Letters*. 2006; 96:208701.
9. Zanette DH. Coevolution of agents and networks in an epidemiological model. 2007 arxiv: 0707.1249.
10. Volz E, Meyers LA. Susceptible-infected recovered epidemics in dynamic contact networks. *Proc Roy Soc B*. 2007; 274:2925–2933.
11. Volz E, Meyers LA. Epidemic thresholds in dynamic contact networks. *J Roy Soc Interface*. 2009; 6:233–241. [PubMed: 18664429]
12. van Segroek S, Santos FC, Pacheco JM. Adaptive contact networks change effective disease infectiousness and dynamics. *PLoS Comp Biol*. 2010; 6(8):e1000895.
13. Gil S, Zanette DH. Coevolution of agents and opinions: Opinion spreading and community disconnection. *Physics Letters A*. 2006; 356:89–94.
14. Zanette DH, Gil S. Opinion spreading and agent segregation on evolving networks. *Physica D*. 2006; 224:156–165.
15. Kozma B, Barrat A. Consensus formation on adaptive networks. *Phys Rev E*. 2008; 77:016102.
16. Kozma B, Barrat A. Consensus formation on coevolving networks: groups' formation and structure. *J Math Phys A*. 2008; 41:224020.
17. Iniguez G, Kertész J, Kaksi KK, Barrio RA. Opinion and community formation in coevolving networks. *Phys Rev E*. 2009; 80:066119.
18. Vazquez F, González-Avella JC, Eguíluz VM, San Miguel M. Time-scale competition leading to fragmentation and recombination transitions in the coevolution of network and states. *Phys Rev E*. 2007; 76:046120.
19. Holme P, Newman MEJ. Nonequilibrium phase transition in the coevolution of networks and opinions. *Phys Rev E*. 2006; 74:056108.
20. Granovetter M. Threshold Models of Collective Behavior. *American Journal of Sociology*. 1978; 83:1420.
21. Watts DJ. A simple model of global cascades on random networks. *Proc Nat'l Acad Sci*. 2002; 99:5766–5771. [PubMed: 16578874]

22. Galstyan A, Cohen P. Cascading dynamics in modular networks. *Phys Rev E*. 2007; 75:036109.
23. De Bruyn A, Lilien GL. A multi-stage model of word-of-mouth influence through viral marketing. *Int J Res Mark*. 2008; 25:151.
24. Friedkin NE. The Attitude-Behavior Linkage in Behavioral Cascades. *Soc Psychol Q*. 2010; 73:196.
25. Melnik S, Ward JA, Gleeson JP, Porter MA. Multi-stage complex contagions. *Chaos*. 2013; 23:013124. [PubMed: 23556961]
26. Durrett, R. *Probability: Theory and Examples*. 4. Cambridge U. Press; 2010.
27. Durrett, R. *Random Graph Dynamics*. Cambridge U. Press; 2008.
28. Kimura D, Hayakawa Y. Coevolutionary networks with homophily and heterophily. *Phys Rev E*. 2008; 78:016103.
29. Vazquez F, Eguíluz VM, San Miguel MS. Generic absorbing transition in coevolution dynamics. *Phys Rev Letters*. 2008; 100:108702.
30. Herrera JL, Cosenza MG, Tucci K, González-Avella JC. General coevolution of topology and dynamics in networks. *Europhysics Letters*. 2011; 95:58006.
31. Böhme GA, Gross T. Analytical calculation of fragmentation transitions in adaptive networks. *Phys Rev E*. 2011; 83:035101.
32. Durrett R, Gleeson J, Lloyd A, Mucha P, Shi F, Sivakoff D, Socolar J, Varghese C. Graph fission in an evolving voter model. *Proc Nat'l Acad Sci*. 2012; 109:3682–3687. [PubMed: 22355142]
33. Böhme GA, Gross T. Fragmentation transitions in multistate voter models. *Phys Rev E*. 2012; 85:066117.
34. See Supplemental Material at [URL will be inserted by publisher] for the finite size scalings of the timescales, the dependence on the average degree, and diagnostics for the numerical fittings.

FIG. 1.

Fraction of the minority opinion at the end as the rewiring probability α varies, starting from Erdős-Rényi graphs with $N=100,000$ nodes and average degree $\lambda = 4$ [32]. Each curve corresponds to a different initial density u of 1's. The final fractions of the minority for α below the $\alpha_c(u)$ phase transitions follow a universal curve independent of the initial fractions. For convenience, the values $\alpha_c(0.5) \doteq 0.73$ and $\alpha_c(0.1) \doteq 0.47$ have been indicated in the figure.

FIG. 2.

Freezing time (i.e., number of updates needed to reach the frozen configuration) at each rewiring probability α . Each point is an average of 20 simulations, starting from Erdős-Rényi graphs with $N = 10^4$ nodes and average degree $\lambda = 4$. Each curve corresponds to a different initial density u of 1's. Below the critical α freezing takes almost $O(N^2)$ updates while it only requires $O(N \log N)$ updates above criticality. Again for convenience, the values $\alpha_c(0.5) \doteq 0.73$ and $\alpha_c(0.1) \doteq 0.47$ have been indicated in the figure.

FIG. 3.

Top: Evolution of the fraction of edges that are discordant 0-1 edges, $N_{01}(t)/M$, versus the fraction of 1's $N_1(t)/N$ when $\alpha = 0.5$ for the rewiring-to-random dynamic. Five simulations starting from $u=0.2, 0.35, 0.5, 0.65$, and 0.8 are plotted in different colors. Each simulation starts from an Erdős-Rényi graph with $N=100,000$ nodes and average degree $\lambda = 4$. After initial transients, the fraction of discordant edges behaves as a function of the population of opinions. Bottom: Similarly, the number of 0-1-0 connected triples behaves as a function of N_1/N after an initial transient (one simulation).

FIG. 4.

Top: plot of the fraction of discordant edges versus the population of opinions in barycentric coordinates for three opinions and $\alpha = 0.5$. Multiple simulations corresponding to different initial densities are shown while each one starts from an Erdős-Rényi graph with $N=10,000$ nodes and average degree $\lambda = 4$. Similar to the two-opinion case, the simulations quickly converge to a parabolic cap of quasi-stationary distributions. Bottom: top view of the parabolic caps of quasi-stationary distributions for $\alpha=0.1, 0.2, \dots, 0.8$. We fit the parabolic cap (2) to simulation data at various α 's and then plot the level sets $E_u N = 0$, which are the intersections of the parabolic caps with the $N = 0$ plane, as the large circles with colors indicating values of α .

FIG. 5.

Fitted coefficient $c_0(\beta)$ (left) and $c_2(\beta)$ (right) in (2) for models with multiple opinions. Each value of the coefficients is obtained by fitting (2) to multiple simulations starting from Erdős-Rényi graphs with $N=100,000$ nodes and average degree 4. The 95% confidence intervals for the 5-opinion fitted coefficients are shown as representative for the other coefficients, but they are barely visible since the standard errors are extremely small (see the Supplemental Materials [34] for complete coefficient fitting results). A linear regression of c_0 on β yields $c_0 = a + b\beta$ where a is insignificant and $b = 0.26 \pm 0.01$ (with 95% confidence bounds). Similarly the linear regression of c_2 on β yields $c_2 = a + b\beta$ where $a = 1.35 \pm 0.03$ and $b = 0.52 \pm 0.01$.

FIG. 6.

Fitted coefficients $c_0(\beta)$ (left) and $c_2(\beta)$ (right) in (4) for models with multiple opinions. Each value of the coefficients is obtained by fitting (4) to the same data as in Figure 5, and the 95% confidence intervals for the coefficients in the 5-opinion model are shown (as representative for the other coefficients, see again the Supplemental Materials [34] for complete coefficient fitting results). A linear regression of c_0 on β yields $c_0 = a + b\beta$ where a is insignificant and $b = 0.22 \pm 0.01$ (with 95% confidence bounds). Similarly the linear regression of c_2 on β yields $c_2 = a + b\beta$ where $a = 1.40 \pm 0.26$ and $b = 0.36 \pm 0.06$.

FIG. 7.

Coefficient $c_3(\beta)$ in (4) for models with multiple opinions. Each value of $c_3(\beta)$ is obtained by fitting (4) to the same data in Figure 5.

FIG. 8.

Number of opinions observed at freezing for different $\beta = \alpha/(1-\alpha)$, starting from k equally likely opinions. Each data point corresponds to a simulation starting from an Erdős-Rényi graph with 100,000 nodes and average degree 4. Results above β_k when starting from k opinions are omitted because the final density of opinions under such conditions is the same as the initial density.