Published in final edited form as:

Technology (Singap World Sci). 2014 March ; 2(1): 36–. doi:10.1142/S2339547814500010.

Diffeomorphometry and geodesic positioning systems for human anatomy

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Abstract

The Computational Anatomy project has largely been a study of large deformations within a Riemannian framework as an efficient point of view for generating metrics between anatomical configurations. This approach turns D'Arcy Thompson's comparative morphology of human biological shape and form into a metrizable space. Since the metric is constructed based on the geodesic length of the flows of diffeomorphisms connecting the forms, we call it *diffeomorphometry*. Just as importantly, since the flows describe algebraic group action on anatomical submanifolds and associated functional measurements, they become the basis for positioning information, which we term *geodesic positioning*. As well the geodesic connections provide Riemannian coordinates for locating forms in the anatomical orbit, which we call *geodesic coordinates*. These three components taken together — the metric, geodesic positioning of information, and geodesic coordinates — we term the *geodesic positioning system*. We illustrate via several examples in human and biological coordinate systems and machine learning of the statistical representation of shape and form.

The low-dimensional matrix Lie groups form the core dogma for the now classic study of the kinematics of rigid bodies in the field of rigid body mechanics. Their infinite dimensional analogue, the diffeomorphism group, containing one-to-one smooth transformations^{1–6}, plays the central role in studying deformable structures in the field of Computational Anatomy (CA)^{7–15}. Central to CA is the comparison of shape and form, morphology, as pioneered by D'Arcy Thompson¹⁶. We are focusing on shapes formed by the submanifolds of the human body in R^3 . Comparison of their coordinates are described via flows of diffeomorphisms connecting them. Thompson's morphological space is made into a metrizable space via a metric induced by the geodesic lengths of the flows. This

reduction of the space of human shape and form to a metric space via diffeomorphic connection we call diffeomorphometry.

We formulate the diffeomorphic correspondence as positioning via the design of optimal control strategies flowing information from one anatomical coordinate system onto another. The flows of the coordinates $t \mapsto \phi_t$ act to transfer the image information, with the controls the vector fields $t \mapsto v_t$ acting as the push. They are related via the dynamics of the system $\dot{\phi} = v \circ \phi$. To define the geodesic shortest path flows a least-action principle is introduced based on a Lagrangian on the generalized coordinates of this system. This links us to classical formulations from physics. The Lagrangian has the property that the metric between any two forms is unchanged by coordinate transformation identically applied to both. Importantly, this property termed right-invariance of the metric, implies that a linear functional of the vector field encoding the geodesic connection between forms satisfies a conservation law.

This linear functional we term the Eulerian momentum. Importantly, conservation implies that the metric comparison of diffeomorphometry is encoded via one of the forms and the Eulerian momentum specifying the geodesic connection of it to the second form. Our parametric reduction to coordinates of the shape or anatomical phenotype is the parametric representation of these linear momenta.

The flows of diffeomorphisms both encode shape as well as transfer the physiological information stored in the imagery. This approach models the observable space of imagery and forms as an orbit under the diffeomorphism group action. We call this orbit the morphological space of forms. It is an example of a Grenander deformable template¹⁷, reducing the study of the image or form to the study of the templates and the transformations that are applied. The transfer of information as group action implies the metric property between coordinate systems is inherited by the elements in the orbit itself. This is vital for the quantitative representation of function across multiscale models which link cells and organ systems that characterize variation in health and disease in Computational Medicine¹⁸.

This diffeomorphic action is a positioning system providing our space of forms with coordinates, termed Riemannian exponential or geodesic coordinates. To make an analogy in these infinite dimensional coordinatized spaces with coordinates on Earth, choosing a template in CA is equivalent to selecting an origin on the globe, say the North pole. Geodesics stemming from the template are analogous to longitude lines, or great circles stemming from the pole, and CA coordinates select both the geodesic and how far to go along it, similar to providing longitude and latitude in spherical coordinate systems.

Noteably, the diffeomorphic actions are rich enough that the morphological space is homogeneous, so that any form can generate the full orbit and can be the center of a chart creating a local coordinate system with geodesic coordinates. The collection of geodesic maps covers the morphological space and defines the geodesic atlas. Geodesic positioning provides a systematic way to choose the chart which flattens to first order the metric, providing exact distances to the origin radially.

Because the image information indexed to submanifolds $M \subset R^3$ are so essential for understanding medicine at the human organ system level of subcortical and surface coordinates in brain and heart, our optimal control strategies are reformulated via Hamiltonian reduction of the Lagrangian on R^3 to the local coordinates $M \subset R^3$ of physiological significance within the body.

METHODS

The morphological space: transformation of forms via diffeomorphisms

CA examines the interplay between imaged anatomical structures indexed with coordinates, hereafter called a *form* lying in some set $\mathcal{M}$ called the *morphological space*, and acted upon by some group of transformations, $\phi \in G$. Our algebraic model of the interaction between the pair (ϕ, m) is that the form $m \in \mathcal{M}$ is carried by the coordinate system represented by $\phi \in G$, and denoted algebraically as group action

$$(\phi, m) \mapsto \phi \cdot m \in \mathcal{M}. \quad (1)$$

The forms studied in CA are submanifolds (points, curves, surfaces, subvolumes) in the human body, and dense scalar and tensor imagery. The transformations are diffeomorphisms $\phi \in G$ a group of one-to-one, smooth, coordinate change of the background space R^3 , with law of composition $\phi \circ \phi'(\cdot) = (\phi(\phi'(\cdot)))$, inverse ϕ^{-1} . In functional anatomy¹⁹, diffeomorphisms represent the structural phenotypes, with imagery functional phenotypes. This separation is not always distinct.

The diffeomorphic transformations are generated as flows. If v_t is a time-dependent vector field on R^3 , the differential equation flow $\dot{y} = v_t(y)$ is given with its inverse, for $t \in [0, 1]$, by

$$\frac{d\phi_t}{dt} = v_t \circ \phi_t, \quad \frac{d(\phi_t^{-1})}{dt} = -d(\phi_t^{-1})v_t, \quad (2)$$

where $df = \left(\frac{\partial f_i}{\partial x_j} \right)$ denotes the 3×3 Jacobian matrix of $f: R^3 \rightarrow R^3$, and both equations are solved with initial condition $\phi_0 = \phi_0^{-1} = \text{id}$, where $\text{id}(x) = x$ is the identity mapping. Here $\dot{\phi} \in R^3$ is the Lagrangian velocity indexed to the initial body configuration, with $v_t \in R^3$ the Eulerian vector field indexed to the flow ϕ_t as a function of time t . The flow of (2) generates a well defined time-dependent path of C^1 diffeomorphisms $t \rightarrow \phi_t$ with inverse (the transport equation) for suitable control of the spatial derivatives of v_t (see below).

Our two kinds of forms are manifolds and imagery, carried by the coordinate transformations of (1). For collections of submanifolds (landmarks, curves, or surfaces) a form M is a subset, usually the geometrical support of a submanifold. Often the manifold of points are parameterized via vertex representations as surfaces in brain and cardiac studies^{20–23}, implying $q(U) = M$ for some mapping $q: U \rightarrow R^3$, where U is taken to be a subset of Euclidean space, although U can be discrete sets or submanifolds. For MRI the

forms are images $I:R^3 \rightarrow R^m$, dense scalar or vector functions, action as right inverse. This gives our two actions:

$$(\varphi, M) \mapsto \varphi \cdot M \doteq \varphi(M), \quad (\varphi, q) \mapsto \varphi \cdot q = \varphi \circ q. \quad (3a)$$

$$(\varphi, I) \mapsto \varphi \cdot I \doteq I \circ \varphi^{-1}, \quad \varphi \cdot I(x) = I(\varphi^{-1}(x)). \quad (3b)$$

There can be many actions as shown in the extended methods for vector imagery and tensor imagery as 3×3 non-negative symmetric matrices, as well as frames or orthonormal bases. Defining in this way specifies the transformation is group action, $\phi, \phi' \in G$, $(\phi \circ \phi') \cdot m = \phi \cdot (\phi' \cdot m) \in \mathcal{M}$.

Diffeomorphometry via geodesic connection

For turning the “ology” of D’Arcy Thompson’s morphology into the “ometry” of diffeomorphometry we construct a metric $\rho_{\mathcal{M}}$ on the morphological space $\mathcal{M}$ of forms, induced by the metric between diffeomorphic changes in coordinates. For this, let $\phi:[0,1] \rightarrow G$ be a differentiable path connecting coordinate systems $\phi_0 = g$, $\phi_1 = h \in G$. The metric between coordinate transformations is defined as the length of the geodesic connection given by the integrated tangent norm along the flow, which induces the metric between forms $m, n \in \mathcal{M}$ defined via the exemplar or template m_0 :

$$\rho(g, h) = \inf_{\varphi: \varphi_0=g, \varphi_1=h} \int_0^1 \|\dot{\varphi}_t\|_{\varphi_t} dt, \quad (4a)$$

$$\rho_{\mathcal{M}}(m, n) = \inf_{g, h: g \cdot m_0=m, h \cdot m_0=n} \rho(g, h). \quad (4b)$$

The norm-square for the tangent is taken as $\|\dot{\varphi}_t\|_{\varphi_t}^2 = \|\dot{\varphi}_t \circ \varphi_t^{-1}\|_V^2 = \|v_t\|_V^2$, where $\|\cdot\|_V$ is the norm for the Hilbert space of vector fields V . To ensure flows satisfying Equation (2) have well defined inverse (requiring Jacobian) and Equation (4) has well defined length for every pair of diffeomorphisms, we constrain the vector fields to be in the Hilbert space $v \in V$ modeled as continuously embedded in continuous, differentiable in space, vector fields with supremum norm so that Eulerian velocities $v \in V$ are spatially C^1 .

Statement 1—Define the group G as^{3,4}

$$G \doteq \left\{ \varphi = \phi_1: \int_0^1 \|v_t\|_V dt < \infty, \dot{\phi}_t = v_t \circ \phi_t, \phi_0 = \text{id} \right\},$$

with orbit of forms $\mathcal{M} = \{m: m = \phi \cdot m_0, \phi \in G\}$.

- Then $\mathcal{M}$ is homogeneous under action G , so that for all $m, m' \in \mathcal{M}$, there exists $\phi \in G$, such that $m' = \phi \cdot m$.

- With $\|\dot{\varphi}_t\|_{\varphi} = \|\dot{\varphi}_t \circ \varphi_t^{-1}\|_V = \|v_t\|_V$ in Equation (4), the metric is right-invariant, for all $\psi \in G$,

$$\rho(g \circ \psi, h \circ \psi) = \rho(g, h). \quad (5)$$

The group action property implies getting from one form to another is possible, extending to morphological spaces the property familiar from the kinematics of rigid bodies in Euclidean space, called homogeneity of the action; see Fig. 1.

The metric being right-invariant implies that (i) the metric on the orbit of forms satisfies the triangle inequality in spite of there being many diffeomorphisms which are indistinguishable seen through their action²⁴, and (ii) computation of geodesics reduces to *shooting* from the identity.

It also gives us a *consistency* property when manipulating the forms. We would prefer our metric to have the property that it obtains consistent measurements across different scales, parameterizations, or extent of inclusion of subparts. For this we define the notion of the *extension* of one form from another. We say $\mathcal{M}'$ is an extension of $\mathcal{M}$ if there exists an onto mapping $\pi: \mathcal{M}' \rightarrow \mathcal{M}$ that is consistent with the group action: $\pi(\phi \cdot m') = \phi \cdot \pi(m')$, $\phi \in G$.

The right invariance of the metric is necessary and sufficient for metric consistency to hold.

Statement 2

- The metric $\rho_{\mathcal{M}}$ on forms $m, n \in \mathcal{M}$ shoots from identity, satisfying triangle inequality²⁴:

$$\rho_{\mathcal{M}}(m, n) = \inf_{g: g \cdot m = n} \rho(\text{id}, g). \quad (6a)$$

- The metric $\rho_{\mathcal{M}}$ on forms is consistent, for any extension $\pi(\phi \cdot m') = \phi \cdot \pi(m')$, then

$$\rho_{\mathcal{M}}(m, n) = \inf_{m': \pi(m') = m, n': \pi(n') = n} \rho_{\mathcal{M}'}(m', n'). \quad (6b)$$

Importantly this provides a straightforward mechanism for a purely geometric comparison or shape registration of manifolds, independent of parametrization by implicit or explicit optimization over all possible reparametrizations, finding the most favorable one for the distance²⁵.

Geodesic coordinates and positioning

We compute geodesics as variational minimizers along the integrated Lagrangian or kinetic

energy $\mathcal{L}(\varphi, \dot{\varphi}) = \frac{1}{2} \|v\|_V^2$; computation of geodesics corresponds to a least-action principle:

$$\rho(g, h)^2 = \inf_{\substack{v: \dot{\varphi}_t = v_t \circ \varphi_t \\ \varphi_0 = g, \varphi_1 = h}} \left(2 \int_0^1 \mathcal{L}(\varphi_t, \dot{\varphi}_t) dt = \int_0^1 \|v_t\|_V^2 dt \right). \quad (7)$$

For computing geodesics we use a matrix differential operator L (powers of the Laplacian) to rewrite the norm in V giving $\|v\|_V^2 = (Lv|v)$ and inner-product for $v, w \in V$, taken componentwise $(Lv|w) \doteq \sum_{i=1}^3 \int_{\mathbb{R}^3} L v_i(x) w_i(x) dx$. We call $\mu = Lv$ the *Eulerian momentum*; sometimes it is a measure.

The Euler equation dictates the temporal evolution of the momentum of the geodesics. The right-invariant metric implies conservation of Eulerian momentum so that the geodesics are encoded by their initial condition, the basis for shooting. In Eulerian coordinates, in passing from $t \mapsto t + \varepsilon$, the particles are at different positions, so for conservation the function being acted upon must be transformed by the adjoint $w \mapsto ((d\varphi_t)w) \circ \varphi_t^{-1}$.

Statement 3—The geodesics minimizing the action of (7) with $\|v\|_V^2 = (Lv|v)$ have Euler momentum satisfying force evolution and conservation on smooth functions $w \in V$:

$$\left(\frac{dLv_t}{dt} | w \right) + (Lv_t | (dv_t)w - (dw)v_t) = 0, \quad (8a)$$

$$(Lv_t | ((d\varphi_t)w) \circ \varphi_t^{-1}) = (Lv_0 | w). \quad (8b)$$

Moreover, geodesics have constant speed $\|v_0\|_V$ and Lagrangian $\mathcal{L}(\varphi, \dot{\varphi}) = \frac{1}{2} \|v_0\|_V^2$ (proven in¹⁰).

The operations consisting in finding an optimizing geodesic in Equations (7) and solving (8) with initial condition Lv_0 are inverse. Under suitable conditions, solutions of the latter are minimizers of the former and conversely. This “inverse” relationship is referred to as the Riemannian exponential and logarithm, with the logarithm only well defined when there are unique solutions to the geodesic equation.

Statement 4—We term *geodesic positioning* as the Riemannian exponential at the identity $\text{Exp}_{\text{id}}(\cdot): V \rightarrow G$ given by the geodesic satisfying Equation (8) with initial condition $v_{t=0} = v_0$; *geodesic coordinates* are the Riemannian logarithm at the identity $\text{Log}_{\text{id}}(\cdot): G \rightarrow V$ (assuming uniqueness) given by $v_{t=0}$ of the field satisfying geodesic connection:

$$\text{Exp}_{\text{id}}(v_0) = g, \quad \text{with } \text{Log}_{\text{id}}(g) = v_0. \quad (9)$$

Extending to the entire group gives

$$\begin{aligned} \text{Exp}_g(v_0 \circ g) &\doteq \text{Exp}_{\text{id}}(v_0) \circ g \quad \text{with} \\ \text{Log}_g(h) &\doteq \text{Log}_{\text{id}}(h \circ g^{-1}) \circ g. \end{aligned} \quad (10)$$

When the logarithm exists we have $\text{Exp}_{\text{id}}(\text{Log}_{\text{id}}(g)) = g$, but the converse equation $\text{Log}_{\text{id}}(\text{Exp}_{\text{id}}(v_0)) = v_0$ may fail to be true, although it holds if v_0 is “small enough”.

As depicted in Fig. 1, the top panel shows diffeomorphic actions which are rich enough that any form can generate the full orbit (homogeneity of the space). The bottom panel shows each form can act as the center of a chart creating a local coordinate system with geodesic coordinates. In this infinite dimensional setting, any one could be selected, analogous to selection of the North pole in finite dimensional spherical representations. The collection of geodesic maps covers the morphological space and defines the geodesic atlas. The metric flattens to the tangent providing exact distances to the origin radially of the great circles stemming from the template given by $\rho_{\mathcal{M}}(m, n) = \|v_0\|_V$. The fact that our forms are defined by group actions implies this is true for all of the actions defined, including imagery and manifolds.

Computation of Geodesic Positioning System (GPS) via optimal control and Hamiltonian reduction

We solve for geodesic position and coordinates for Log and Exp as an inexact matching of coordinate systems using classical methods in optimal control. In CA, vector fields are used as control that push the forms (submanifolds or images) within the anatomical morphological space. We take as the dynamical system the flow of coordinates $t \mapsto q_t \doteq \phi_t \cdot q_0 \in Q$ modelled as embedded in some linear state space $Q \supset M$, and the control the vector field $t \mapsto v_t \in V$ related to the state via the infinitesimal action of the flow $\dot{q}_t = v_t \cdot q_t, q_{t=0} = q_0$. The optimal control takes the running cost the kinetic energy $(Lv_t|v_t) = \|v_t\|_V^2$, with an assumed target or endpoint condition E.

Control Problem 1

$$\min_{v: \dot{\varphi}_t = v_t \circ \varphi_t, \varphi_0 = \text{id}} C(v) \doteq \frac{1}{2} \int \|v_t\|_V^2 dt + E(\varphi_1 \cdot q_0) \quad (11a)$$

$$\min_{v_0} C(v_0) \doteq \frac{1}{2} \|v_0\|_V^2 + E(\text{Exp}_{\text{id}}(v_0) \cdot q_0). \quad (11b)$$

The control indexed to initial v_0 is *shooting* and has been examined in various forms^{4,10,25–32}. Various endpoint conditions have been examined, including for dense image

and tensor matching $E = \frac{1}{2} \|\varphi \cdot I - J\|_2^2$, for point matching with correspondence

$E = \frac{1}{2} \sum_i |y_i - \varphi(x_i)|^2$, as well as other for point-sets, curves and surfaces without correspondence.

Forms in the orbit being positioned can be of reduced dimension to the diffeomorphisms in the group. Manifolds represented by multiple organs or fiducial points are singular in the background space indexing the vector fields. The GPS computation exploits this by switching to a Hamiltonian formalism in which the momentum is reparameterized in terms of a co-state variable p indexed to the lower complexity state space of the forms. Introduce a Lagrange multiplier p_t dual to the state space for constraining $q_t \dot{=} v_t \cdot q_t$ added to the negative Lagrangian, giving control-dependent Hamiltonian

$$H_v(q, p) = (p|v \cdot q) - \frac{1}{2}(Lv|v). \quad (12)$$

The Pontryagin Maximum principle yields the optimized geodesic control $\hat{v}$, and the Hamiltonian $H(q, p) = \max_v H_v(q, p)$ with dynamics

$$\begin{pmatrix} \dot{q} \\ \dot{p} \end{pmatrix} = \begin{pmatrix} \partial H(q, p) / \partial p \\ -\partial H(q, p) / \partial q \end{pmatrix} \doteq F(q, p), \quad q_0 = \text{fixed}. \quad (13)$$

Our *shooting* problem for $q_t = \phi_t \cdot q_0$, $q_t \dot{=} v_t \cdot q_t$ has been reduced to the geodesic control $\hat{v}$ with momentum Lv satisfying Euler Equation (8) with Hamiltonian constant equalling the Lagrangian:

$$H(q, p) = H(q_0, p_0) = \frac{1}{2} \|\hat{v}_0\|_v^2.$$

The reduced control problem becomes as follows.

Control Problem 2—Hamiltonian Reduction

$$\min_{p_0} C(p_0) = H(q_0, p_0) + E(q_1) \text{ subject to} \quad (14)$$

$$\begin{pmatrix} \dot{q} \\ \dot{p} \end{pmatrix} = F(q, p), \quad q_0 = \text{fixed}.$$

For manifolds $M = q(U)$, $q : U \rightarrow R^3$, $q_t = \phi_t \circ q_0$, the infinitesimal action is $v \cdot q \doteq v \circ q$ with Lagrange multiplier co-state constraint

$$v \cdot q \doteq v \circ q, \quad (p|v \circ q) \doteq \int_U \langle p(u), v \circ q(u) \rangle du. \quad (15)$$

Maximizing $H_v(p, q)$ of (12) in v with $K = L^{-1}$ gives Dirac geodesic momentum and optimizing vector field satisfying Equation (8):

$$L\hat{v} = \int_U \delta_{q(u)} \otimes p(u) du, \quad \hat{v}(\cdot) = \int_U K(\cdot, q(u)) p(u) du. \quad (16)$$

The manifold dynamical system becomes

$$\begin{pmatrix} \dot{q} \\ \dot{p} \end{pmatrix} = \begin{pmatrix} \hat{v} \circ q \\ -(d\hat{v})^* \circ qp \end{pmatrix} = F(q, p), \quad (17a)$$

$$\begin{aligned} H(q, p) &= \frac{1}{2} \int_{U^2} p^*(u) K(q(u), q(u')) p(u') du' du \\ &= H(q_0, p_0) \doteq \frac{1}{2} (p_0 | K_{q_0 q_0} p_0) = \frac{1}{2} \|\hat{v}_0\|_{\mathcal{V}}^2, \end{aligned} \quad (17b)$$

where we use $(\cdot)^*$ to denote matrix transpose.

For dense imagery, $q_t = I \circ \varphi_t^{-1}$, $q_0 = I$, the infinitesimal action is $v \cdot q = -\langle \nabla q, v \rangle$ with Lagrange multiplier co-state constraint

$$(p | -\langle \nabla q, v \rangle) = - \int_{R^3} p(x) \langle \nabla q(x), v(x) \rangle dx. \quad (18)$$

Maximizing $H_v(p, q)$ of (12) in v gives geodesic momentum and vector field satisfying Equation (8):

$$L\hat{v} = -p\nabla q, \quad \hat{v}(\cdot) = - \int_{R^3} K(\cdot, x) \nabla q(x) p(x) dx. \quad (19)$$

The dense image dynamical system becomes

$$\begin{pmatrix} \dot{q} \\ \dot{p} \end{pmatrix} = \begin{pmatrix} -\langle \nabla q, \hat{v} \rangle \\ \nabla \cdot (p\hat{v}) \end{pmatrix} = F(q, p); \quad (20a)$$

$$\begin{aligned} H(q, p) &= \frac{1}{2} \int_{R^3 \times R^3} (p\nabla q)^*(x) K(x, y) (p\nabla q)(y) dx dy \\ &= H(q_0, p_0) \doteq \frac{1}{2} (p_0 | K_{q_0 q_0} p_0) = \frac{1}{2} \|\hat{v}_0\|_{\mathcal{V}}^2. \end{aligned} \quad (20b)$$

The Extended Methods does the case of atlases as collection of multiple manifolds.

Adjoint method of solution

For solving control problem 2 we use gradient based adjoint methods^{27,33–35} transporting the endpoint condition backwards, reoptimizing with respect to p_0 with Hamiltonian dynamics $F(q, p)$.

The adjoint method arises by introducing the functional determining our gradient method

$$\int_0^1 \left(\lambda \left| \begin{pmatrix} \dot{q} \\ \dot{p} \end{pmatrix} - F(q, p) \right| \right) dt + H(q_0, p_0) + E(q_1), \quad (21)$$

where λ_t is the Lagrange multiplier with

$$\left(\lambda_t \left| \begin{pmatrix} f \\ g \end{pmatrix} \right. \right) = \int_U (\lambda_{t,q}(u)f(u) + \lambda_{t,p}(u)g(u))du.$$

The gradient variation under a perturbation $(q, p) \rightarrow (q, p) + \delta(q, p)$ gives equations for λ :

$$(\dot{\lambda} + dF^*(q, p)\lambda | \delta(q, p)) = 0 \quad (22a)$$

$$(dH(q_0, p_0) - \lambda_0 | \delta(q_0, p_0)) = 0 \quad (22b)$$

$$\lambda_{1,p} = 0, \quad \left(\frac{\partial E(q_1)}{\partial q} + \lambda_{1,q} | \delta q_1 \right) = 0, \quad (22c)$$

where dH is the differential with respect to (q, p) .

See the Extended Methods gradient algorithm Equations (32) for solving the maximizer conditions.

RESULTS

Geodesic positioning in high field MRI

Figure 2 illustrates geodesic positioning for segmentation via deformable templates^{36,37}. The top row shows a high field 11.7T MRI reconstructed hippocampus³⁸ from a normal control which has been partitioned into four substructures, with dense subvolumes and bounding surfaces: M_0 equal to CA1 (blue), CA2 (green), CA3/Dentate gyrus (red), and subiculum (cyan). Taking this as the template M_0 is used to parcellate or segment the target in the lower panel by positioning it onto the target coordinates $\phi \cdot M_0 = \text{Exp}_{\text{id}}(v_0) \cdot M_0$ with $v_0 = \text{Log}_{\text{id}}(\phi)$. The target has been selected manifesting temporal lobe atrophy. The left part of the bottom panel shows the cross sections of the mapped template in the MRI. The right of the bottom panel shows the mapped template with two of the substructures CA1 and CA3/Dentate shown in the atrophied coordinates of the target. These two substructures illustrate the greatest atrophy, with CA3/Dentate gyrus showing the greatest amount of atrophy as much as 25 percent in target coordinates as measured by the determinant of the Jacobian. The geodesic positioning encoded by $\phi = \text{Exp}_{\text{id}}(v_0)$ was determined by associating landmarks to each of the M_0, M_1 structures and solving the optimal Control Prob. 3 via landmark matching¹⁰.

Geodesic positioning of brain DTI tracts

Shown in Fig. 3 is geodesic positioning used by many groups in Computational Anatomy registering information in diffusion tensor MRI (DT-MRI). Top row of Fig. 3 shows submanifold curves generated via FACT tract tracing³⁹ from twenty brains following rigid alignment to MNI space. The commissural fiber connecting the postcentral gyri (CC-PoCG, Fig. 3 left column) and the left cortical-spinal tract (CST-Left, Fig. 3 right column) are shown for demonstration. Following⁴⁰, the MRI coordinates were repositioned aligning the

brains to a common template coordinate system $I_0 \xrightarrow{\varphi^{(i)}} I_i$ solving control problem 1 based on the fractional anisotropy and B0 contrasts in the DWI imagery using Multi-channel LDDMM⁴¹. The bottom row shows the fiber tracts regenerated after aligning the brains to the common template coordinate system.

Exponential geodesic coordinates in brain and machine learning

Figure 4 depicts the use of geodesic coordinates to visualize populations of structures around the templates. Shown is the geodesic positioning of the single hippocampus subcortical structure from a population in the BIOCARD study. For the graph the population was mapped solving control problem 2 to a common template form using template estimation as in^{35,42}. For each structure M_i the diffeomorphism was calculated $M_0 \xrightarrow{\varphi^{(i)}} M_i$, solving Control Prob. 2 for surface matching 27, from which exponential coordinates were calculated $\text{Log}_{\text{id}}(\phi^{(i)}) = v^{(i)}$. Figure 4 shows the two-dimensional PCA representation of the exponential vector field coordinates of the mapped structures. Shown at the center is the template.

Figure 5 is an illustration of using these exponential coordinates for machine learning of temporal lobe structures, amygdala-entorhinal cortex-hippocampus in the BIOCARD study. Shown in the Figure are results of machine learning using linear discriminant analysis (LDA) on the exponential coordinates, with 50% of the data withheld for training. The population was mapped solving control problem 2 to a common template as above and for each structure M_i the diffeomorphism was calculated, from which exponential coordinates were calculated. A finite dimensional basis was generated from the coordinates for each mapping $\{v^{(i)}\}$, expanding in a PCA basis based on the empirical covariance. Up to 50 anatomies are shown placed in the first two LDA coordinates for each group showing red for normal, green for clinical Alzheimer's disease, and blue for those diagnosed subsequent to their last scan and termed pre-clinical. Three of the 50 examples from each group are plotted showing the surface temporal lobe structures.

Geodesic coordinates for heart shape

The geodesic exponential coordinates are the basis for many machine learning and statistical methods for discrimination and indexing of shapes. Figure 6 shows results in hearts. The top shows a reconstruction of a high resolution computed tomography left ventricular template (1 mm isotropic) generated from a population of 25 subjects. Shown in color is the 17 AHA partition⁴³ with each segment represented by one color and using the ontology definitions⁴⁴. The population of 25 heart geometries were positioned solving control problem 2 using LDDMM⁴⁵ for the dense imagery. Figure 6 depicts the use of Riemannian coordinates for statistical encoding of the shape phenotype. For this the population was mapped solving control problem 1 generating coordinates $v^{(i)} = \text{Log}_{\text{id}}(\phi^{(i)})$ defined by template to target maps $I_0 \xrightarrow{\varphi^{(i)}} I_i$ representing the shape phenotype positioned relative the template heart coordinate system. Shown depicted via the black area is the indicated region that was statistically significant between the two different populations of ischemic versus non-ischemic cardiomyopathy at end-systole⁴⁶. The region is AHA segment 13 (anterior apical) where the majority of ischemic patients had their infarction. Bottom panel illustrates the

spatial distribution of average determinant of Jacobian for the ischemic (left) and non-ischemic (right) populations within AHA segment 13. Ischemic patients have smaller Jacobian determinant indicating wall thinning due to scar formation.

DISCUSSION

The fact that the right-invariant diffeomorphometry metric is consistent provides a straightforward mechanism for a purely geometric comparison of manifolds, independent of the parametrization. This can be contrasted with other methods in CA that equip shapes with a predefined parametrization such as Euclidean or spherical coordinates^{47–49}. In the extended methods we illustrate via various reparameterization examples.

It is informative to compare control problem 2 for dense matching of I to J to LDDMM⁵⁰ using Equation (19) relating $Lv \hat{=} -p\nabla q$ with the state $q = I \circ \phi^{-1}$. Dense image matching minimizes endpoint

$$E(q_1) = \frac{1}{2} \|q_1 - J\|_2^2, \quad (23)$$

with $\|\cdot\|_2$ the squared-error norm over the image. In this case the Hamiltonian provides geodesic coordinates giving the reduction from $Lv_0: R^3 \rightarrow R^3$ to $p_0: R^3 \rightarrow R$, reducing the dimension of $\phi, \dot{\phi} = v \circ \phi$ which are 3-vectors to the reduced dimension co-state scalar field. To calculate p_0 , we use conservation to transport the endpoint condition on p_1 back to the origin giving p_0 . The endpoint for p_1 maximizes control problem 2:

$$p_1 = -\frac{\partial E}{\partial q}(q_1) = -(q_1 - J) = -(I \circ \phi_1^{-1} - J). \quad (24)$$

For $E(q_1)$ smooth, then Lv has vector density $(Lv_t | w) = \int_{R^3} \langle m_t(x), w(x) \rangle dx$ with conservation Equation (8b) obtaining the initial density,

$$m_0 = |d\phi_1| (d\phi_1)^* m_1 \circ \phi_1. \quad (25)$$

Using $q_0 = q_1 \circ \phi_1$, $\nabla q_0 = (d\phi_1)^* \nabla q_1$ and Equations (19) and (24), gives maximizer conditions of LDDMM⁵⁰:

$$\begin{aligned} L\hat{v}_1 = m_1 &= -p_1 \nabla q_1 = -(I \circ \phi_1^{-1} - J) \nabla (I \circ \phi_1^{-1}). \\ L\hat{v}_0 = m_0 &= |d\phi_1| (d\phi_1)^* m_1 \circ \phi_1 = |d\phi_1| (J \circ \phi_1 - I) \nabla I. \end{aligned} \quad (26)$$

The geodesic trajectory in the group is characterized by the initial Eulerian momentum from Equation (19) given by $Lv_0 \hat{=} -p_0 \nabla q_0$, $q_0 = \nabla I$ fixed for the template, with $p_0: R^3 \rightarrow R$ the reduced co-state,

$$p_0 = -|d\phi_1| (I - J \circ \phi_1). \quad (27)$$

For pointset correspondence matching $\{x_i\}$ to $\{y_i\}$, then $U = \{x_i\}$ and state $q_i = \phi(x_i)$ with smooth endpoint and vector endpoint condition

$$E = \frac{1}{2} \sum_t |y_i - \phi(x_i)|^2, \quad p_1 = -\frac{\partial E}{\partial q}(q_1) = (q_1 - y). \quad (28)$$

Then Equation (16) for Eulerian momentum is singular,

$$L\hat{v}_t = \sum_i \delta_{q_{ti}} \otimes p_{ti}, \quad \hat{v}_t(\cdot) = \sum_i K(\cdot, q_{ti}) p_{ti}. \quad (29)$$

The conservation Equation (8b) for pointsets gives $p_{0i} = (dq_{1i})^* p_{1i}$, implying the reduction

$$p_{0i} = (dq_{1i})^* (q_{1i} - y_i), \quad L\hat{v}_0 = \sum_i \delta_{q_{0i}} \otimes (dq_{1i})^* (q_{1i} - y_i).$$

EXTENDED METHODS

Group actions

Here are several other group actions. For images represented by vector fields, $I: R^3 \rightarrow R^3$, then

$$\varphi \cdot I = ((d\varphi)I) \circ \varphi^{-1} \quad \text{or} \quad \varphi \cdot I = ((d\varphi^{-1})^* I) \circ \varphi^{-1},$$

where $(d\varphi^{-1})^*$ holds for $(d\varphi^*)^{-1}$.

When $I = (I_1, I_2, I_3)$ is a vector field of frames, specifically a positively oriented orthonormal basis of R^3 denoting the tangent to some curves, and the two other normal vectors forming the Frenet frame, then

$$\varphi \cdot I = \left(\frac{(d\varphi)I_1}{\|(d\varphi)I_1\|}, \frac{(d\varphi^{-1})^* I_3 \times (d\varphi)I_1}{\|(d\varphi^{-1})^* I_3 \times (d\varphi)I_1\|}, \frac{(d\varphi^{-1})^* I_3}{\|(d\varphi^{-1})^* I_3\|} \right) \circ \varphi^{-1},$$

where we have denoted vector cross product as $I_i \times I_j$. The interpretation is that I_1 deforms as a fiber (a tangent to some curve), I_3 deforms like a normal to the plane generated by I_1, I_2 and the deformation of I_2 is uniquely constrained by the fact that the basis is positive and orthonormal.

For tensor images 3×3 non-negative symmetric matrices, Alexander and Gee⁵¹ used rotation of the eigenfunctions via the previous action on frames with eigenvalues unchanged; a second action becomes

$$\varphi \cdot I \doteq ((d\varphi)I(d\varphi)^*) \circ \varphi^{-1}.$$

Consistency of metric

To understand geometric comparison of manifolds, independent of the parametrization examine $q' : U' \rightarrow R^3 \in \mathcal{M}'$ as a parametrized surface with $\pi(q) \in \mathcal{M}$ the restriction of q to a finite set $U \subset U'$ of fiducial points. The statement 2 asserts that the construction that projects $\rho_{\mathcal{M}'}$ onto $\rho_{\mathcal{M}}$ via the right-hand side of (6) is equivalent to applying the construction in statement 2 directly on $\mathcal{M}$. Another important consequence is that changing the parametrization from $U \rightarrow R^3$ to $\tilde{U} \rightarrow R^3$ when U and $\tilde{U}$ are diffeomorphic does not affect the metric.

This is again a purely geometric comparison of manifolds, independent of the parametrization. Take $\mathcal{M}$ to be the set of mappings $q : U \rightarrow R^3$ such that $q(U)$ is a submanifold of R^3 (i.e., the set of embeddings) and $\mathcal{M}$ the set of submanifolds of R^3 . Define $\pi(q) = q(U)$ which associates to each mapping q the subset of R^3 that it parametrizes. Then, $\phi \circ \pi(q) = \phi \circ q(U) = \pi(\phi \circ U)$ and statement 2 implies a distance between submanifolds is obtained

$$\rho_{\mathcal{M}}(m, n) = \inf_{q, \tilde{q} : q(U)=m, \tilde{q}(U)=n} \rho_{\mathcal{M}'}(q, \tilde{q}).$$

Atlases of multiple submanifolds

This generalizes for atlases collections of manifolds $q = (q^1, \dots, q^n)$, $q^i : U_i \rightarrow R^3$, with infinitesimal action $v \cdot q \doteq (v \circ q^1, \dots, v \circ q^n)$. Define $p = (p^1, \dots, p^n)$ yielding

$$(p|v \cdot q) = \sum_{i=1}^n (p^i|v \circ q^i), \text{ optimal control}$$

$$\hat{v}(\cdot) = \sum_i \int_{U_i} K(\cdot, q^i(u)) p^i(u) du. \quad (30)$$

The dynamics and Hamiltonian becomes

$$\begin{pmatrix} \dot{q}^i \\ \dot{p}^i \end{pmatrix} = \begin{pmatrix} \hat{v} \circ q^i \\ -(d\hat{v})^* \circ q^i p^i \end{pmatrix} = F^i(q, p) \quad (31a)$$

$$H(q, p) = H(q_0, p_0) = \frac{1}{2} \sum_{i,j} (p_0^i | K_{q_0^i q_0^j} p_0^j). \quad (31b)$$

Adjoint gradient algorithm

The adjoint algorithm arises by implementing the gradient perturbation $(q, p) \rightarrow (q, p) + \delta q$, p of Equations (21) and (22) to satisfy the fixed point optimizer.

Algorithm 1—Generic Adjoint Initialize p_0 .

i. Given p_0 solve $(q, \dot{p})^* = F(q, p)$, giving p_1, q_1 .

ii. Given q_1 , solve boundary terms $\lambda_{1,q} = -\frac{\partial E}{\partial q}(q_1)$.

iii. Backsolve adjoint equation for λ of (22a):

$$\dot{\lambda}_t = -dF^*(q_t, p_t)\lambda_t. \quad (32a)$$

iv. Compute gradient term (22b) with q_0 fixed and $H(q_0, p_0) = \frac{1}{2}(p_0 | K_{q_0 q_0} p_0)$, giving update

$$p_0^{\text{new}} \leftarrow p_0^{\text{old}} + \varepsilon (K_{q_0 q_0}^{-1} \lambda_{0,p} - p_0^{\text{old}}) \quad (32b)$$

with $\varepsilon = \text{step-size}$ and set $p_0 = p_0^{\text{new}}$ and go to (i).

Acknowledgments

We dedicate this paper to Ulf Grenander the father of Metric Pattern Theory. We thank Drs. Truncoso, Aggarwal and Mori for the 11.7T hippocampus supported by NIH grant EB003543. The authors were supported by grants R01 MH056584, R24 HL085343, R01 EB000975, P50 MH071616, R01 EB001838, P41 EB015909, P41 EB015909, R01 EB008171, R01 MH084803, RR025053-01, U01 NS082085-01 and the ANR HMTc. Data used in Figures 4,5 were obtained from the Bio-markers of Cognitive Decline Among Normal Individuals (BIOCARD) study supported in part by grants from the National Institutes of Health U01-AG033655; a list of Cores can be found at <http://www.biocard-se.org/public/Core%20Groups.html>. The data in Figure 6 was obtained from the Donald W. Reynolds Cardiovascular Clinical Research Center at The Johns Hopkins Hospital supported in part by grants from Reynolds foundation.

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Figure 1.

Top panel shows the orbit of forms $\mathcal{M}$, and acting group G . Bottom panel shows radial representation of geodesic flattening of forms. Tangent space norm and geodesic distance agree for radial great circles emanating from the template m_0 with distance $\rho_{\mathcal{M}}(m_0, n) = \|v_0\|_V$.

Figure 2.

Showing mapping in high field 11.7T hippocampus depicting the partition into CA1 (blue), CA2 (green), CA3/Dentate (red) and Subiculum (cyan) of target structures exhibiting temporal lobe atrophy. Top shows four structures in section from the template corresponding to an age controlled normal showing reconstructions of CA1, CA2, CA3/ Dentate, subiculum. Bottom shows template mapped to the target an individual suffering from temporal lobe atrophy; right structures are CA1, CA3/Dentate colored with Jacobian determinant. Data collected is three-dimensional diffusion tensor imaging and was performed using a horizontal-bore 11.7 T NMR scanner (Bruker Biospin, Billerica, MA). DTI data were acquired with a 3D diffusion-weighted EPI sequence (TE = 27 ms, TR = 500 ms). The imaging field-of-view was 42 mm × 45 mm × 64 mm. Two b0 images and 30 diffusion directions were acquired within a scan time of 13.5 hours for DTI mapping.

Figure 3.

Geodesic positioning of the DT-MRI and alignment of fiber tracts. Three-dimensional diffusion tensor imaging of 20 subjects were performed on a 1.5T Siemens MR unit using single-shot echo-planar imaging sequences with sensitivity encoding (SENSE EPI) with parallel imaging factor of 2.0 (imaging matrix: 96×96 , field-of-view: $240 \text{ mm} \times 240 \text{ mm}$, and slice thickness 2.5 mm). B0 images and DWIs of 30 diffusion directions were acquired and co-registered to remove eddy current and motion. The scanning time was 4 min per dataset. Tensor calculation was performed, followed by rigid alignment of all subjects. Fiber reconstructions were performed in each subject's space using FACT tract tracing algorithm³⁹. CC-PoCG fiber: the starting and ending points of tract tracing were selected as the post-central gyri (PoCG) of the two hemispheres, and the fiber path was constraint to penetrate the corpus callosum (CC). CST_left fiber: the starting and ending points were the cerebral peduncle (CP) and the pre-central gyrus (PrCG); the fiber path was constrained by the posterior limb of internal capsule (PLIC) and the superior corona radiata (SCR). Top row shows the tracts in native brain (4 subjects were shown for demonstration); bottom row shows tracts after geodesic positioning via LDDMM solution of control problem 1.

Figure 4.

Panel shows structures indexed by two highest variance dimensions from PCA on Riemannian exponential coordinates representing hippocampus structures from the BIOCARD study. Each structure is placed on the 2D plane with the template (red) at (0,0) in the center, left/right corresponding to the first dimension, and up/down corresponding to the second.

Figure 5.

Panel shows results of machine learning (LDA) on the geodesic coordinates for sets of temporal lobe structures, amygdala-entorhinal cortex-hippocampus, with 50% withheld for training. Up to 50 points are placed for each group, with three examples from each group shown as surfaces. Red: normal elderly subjects, green: subjects diagnosed with Alzheimer's disease at time of their last scan, blue: subjects diagnosed subsequent to their last scan and termed pre-clinical.

Figure 6.

Top section: A high resolution computed tomography left ventricular template (1 mm isotropic) constructed from 25 subjects in AHA segmentation represented by one color per segment; black area within AHA anterior apical segment 13, showing statistical significance between two different populations of cardiac disease, ischemic ($n = 13$, 10 men, mean age 56) and non-ischemic cardiomyopathy ($n = 12$, 8 men, mean age 52) at end-systole. Each subject was studied either in a 32 ($n = 8$) or 64-detector ($n = 17$) multi-detector computed tomography scanner (Aquilion 32(64), Toshiba Medical Systems Corporation, Otawara, Japan). Plane resolution varied from 0.36×0.36 mm to 0.45×0.45 mm, thickness = 0.5 mm. Bottom section shows average Jacobian for ischemic (left) and non-ischemic (right) groups within segment 13 highlighting that on average non-ischemic group had significantly larger tissue (myocardial) volume relative to ischemic group at end-systole. This indicates smaller wall thickening during maximum contraction at the location of infarction in ischemic population⁴⁶.