Subthalamic–pallidal interactions are critical in determining normal and abnormal functioning of the basal ganglia

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The subthalamic nucleus (STN) and external globus pallidus (GP) form a recurrent excitatory–inhibitory interaction within the basal ganglia. Through a computational model of these interactions we show that, under the influence of appropriate external input, the two nuclei can be switched between states of high and low activity or can generate oscillations consisting of bursts of high-frequency activity repeated at a low rate. It is further demonstrated from the model that the generation of the repetitive bursting behaviour is favoured by increased inhibition of the GP, which is a condition indicated in Parkinson's disease. Paradoxically, increased striatal inhibition of the GP is predicted to cause an increase rather than a decrease in its mean firing rate. These behaviours, arising from a biologically inspired computational model of the STN–GP interaction, have important consequences for basal ganglia function and dysfunction.

Keywords: basal ganglia; subthalamic nucleus; globus pallidus; bifurcation analysis; oscillatory behaviour

1. INTRODUCTION

The basal ganglia (figure 1*a*) have been implicated in movement control and motor planning (Houk & Wise 1995), motor and cognitive sequence generation (Suri & Schultz 1998), sequence encoding (Beiser & Houk 1998), context detection via reinforcement learning (Houk & Wise 1995; Dominey *et al.* 1995) and action selection (Redgrave *et al.* 1999). There is growing realization that processing within the basal ganglia is more complex than is implied in the classical view of a direct and an indirect pathway involving neocortex, basal ganglia and thalamus (Alexander *et al.* 1990; Smith *et al.* 1998). In particular, contemporary neuroanatomy of the basal ganglia reveals a prominent feedback system, involving the excitatory subthalamic nucleus (STN) and the inhibitory globus pallidus (GP) (Smith *et al.* 1994; Shink *et al.* 1996; Joel & Weiner 1997; Wichmann & DeLong 1999). STN projection neurons use glutamate as their primary neurotransmitter and exert a powerful excitatory influence on their targets (Kita 1992). GP projection neurons use GABA (γ -aminobutyric acid) as their primary neurotransmitter and exert an inhibitory influence on their targets (Smith & Bolam 1989). The STN is also subject to excitatory influences from cortex and from thalamus, whilst the GP also receives inhibition from striatum (see figure 1*a*). In addition, local collaterals, suggesting significant interconnectivity within each nucleus, have been described (Iwahori 1978; Kita *et al.* 1983*a*; Kita & Kitai 1994).

Recent analysis (Gillies & Willshaw 1998) has suggested that the presence of even weak interconnectivity will cause the STN to exhibit bistable behaviour. This means that the entire system can reside in a state in which all neurons

fire at close to maximum rate, close to 500 Hz, or in a state with a much lower firing rate. Short-term excitatory input can shift the system into the high state and it was suggested that the GP could provide inhibition to return it to the low state. This implies that large regions of the STN can become uniformly and synchronously active as observed, for example, in rat STN after localized cortical stimulation (Fujimoto & Kita 1993).

In culture preparations of the basal ganglia, Plenz & Kitai (1999) have recently reported electrical oscillations. The generation of the oscillations depends on the interaction between the GP and the STN which mirrors that observed in anatomical studies (Joel & Weiner 1997; Plenz *et al.* 1998). Rhythmic bursting oscillations have also been observed in STN neurons of MPTP (1-methyl-4-phenyl-1,2,3,6-tetrahydropyridine treated) monkeys (Nini *et al.* 1995) and in human patients with Parkinson's disease (Levy *et al.* 2000; Magariños-Ascone *et al.* 2000). Some of the STN cells in these patients have a frequency matching that observed in tremor (5–6 Hz) and are influenced by tremor changes (Magariños-Ascone *et al.* 2000). Other STN cells oscillate at higher frequencies (*ca.* 15 Hz) with evidence of phase locking across the nucleus.

Understanding the underlying dynamics between the STN and GP is essential to understanding these behaviours and how they may be modified or controlled. We have analysed the simplest representative circuit made up of an ensemble of excitatory STN neurons and an ensemble of inhibitory GP neurons (figure 1*b*). The dynamic behaviours observed in this simple system are essentially similar to those in larger-scale systems. In the model, the STN neurons excite the GP and are subject to excitation through local feedback and from cortex. Conversely, the GP neurons inhibit the STN and are subject to inhibition through local feedback and from the striatum. Analysis of this model has enabled us to identify the types of behaviour that this coupled STN–GP dynamical system can

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Figure 1. (a) The major pathways through the basal ganglia. The shaded area denotes the indirect pathway from the striatum to the substantia nigra and entopeduncular nucleus. Blue pathways indicate inhibitory connections, green pathways indicate excitatory connections. The dotted line represents a dopaminergic projection. (b) Schematic diagram of the interactions between and within the STN and GP. x and y indicate the potentials in mV of STN and GP neurons, respectively. a , b , c and d are the coupling parameters for the interactions within and between nuclei, converting firing rates into potentials in mV. For details of the equations used see Appendix A. Abbreviations: CTX, cortex; STR, striatum; STN, subthalamic nucleus; GP, globus pallidus; THAL, thalamus; SNr, substantia nigra pars reticulata; SNc, substantia nigra pars compacta; ENT, entopeduncular nucleus.

Table 1. Parameter values used in the model.

sources	name	value
STN		
Kita <i>et al.</i> (1983b); Afsharpour (1985)	τ^{STN}	6 ms
	$\sigma_{\max}$	500 Hz
Nakanishi <i>et al.</i> (1987)	x_{th}	15 mV
	κ	0.3 mV^{-1}
GP		
Kita & Kitai (1991)	τ^{GP}	14 ms
Nambu & Llinás (1994)	$\xi_{\max}$	100 Hz
	y_{th}	10 mV
	η	0.2 mV^{-1}

generate; in particular the conditions under which the system generates repetitive bursting.

2. MATERIAL AND METHODS

The STN–GP interaction is computed from coupled differential equations describing the mean membrane potentials x , y of populations of STN and GP neurons, respectively, as follows:

$$\begin{aligned} \tau^{STN} dx/dt &= -x + a\sigma(x) - c\xi(y) + I^{CTX} \\ \tau^{GP} dy/dt &= -y - b\xi(y) + d\sigma(x) + I^{STR}, \end{aligned} \quad (2.1)$$

where τ^{STN} and τ^{GP} are the time constants of the STN and GP neurons, respectively. I^{CTX} and I^{STR} are external inputs from the cortex and striatum respectively. a , b , c and d specify connection strengths as indicated in figure 1b. $\sigma(x)$ and $\sigma(y)$ are sigmoidal functions relating the STN and GP potentials x , y to their respective firing frequencies.

We present a mathematical analysis of this system along with accompanying simulation results to demonstrate the behaviours identified. For a general analysis of an excitatory–inhibitory oscillator system see Ermentrout & Cowan (1979). In our modified system, the equations are formulated in terms of membrane

potential to allow direct application of experimental data (e.g. derived voltage firing-rate curves, where the parameters specifying the shapes of the curves are derived from experimental data; see table 1).

Our system is analysed and discussed within the context of STN–GP interactions and parameters. As the strengths of the external inputs and connections between and within the STN and GP are substantially unconstrained, we explore how different strengths of neural connections can influence the types of behaviour the system may exhibit under different strengths of external inputs from cortex and from striatum. The external inputs are assumed to be constant for the period of analysis and we assume that connection delays are negligible. We focus on the general physiological properties that reveal the most salient system behaviours. The model description and assumptions are at the level of general cell properties and the influence of individual channel types cannot be analysed within this system. The behaviour of the model is characterized in terms of the potentials (x_0, y_0) at the equilibrium point, where $dx/dt = dy/dt = 0$, and the dynamic deviations from it. For more mathematical details see Appendix A.

3. RESULTS

(a) Setting of parameter values

The parameters $\sigma_{\max}$, κ , x_{th} in $\sigma(x)$, $\xi_{\max}$, η , y_{th} in $\xi(y)$, τ^{STN} and τ^{GP} are assigned values obtained from the experimental literature as given in table 1. We now describe the effects of varying the parameter values that remain unspecified, namely a , b , c , d , I^{CTX} and I^{STR} . We summarize the mathematical and numerical results that shed light on the following issues: (i) the function of the interconnectivity within the STN; (ii) the effects of cortical and striatal input; (iii) the conditions under which the system is bistable; (iv) the conditions under which oscillations occur; (v) what determines the frequency of

Figure 2. Example simulations of the STN–GP system. In all graphs blue and green lines represent firing rate of the STN and GP, respectively. (a) Bistable behaviour. Brief stimulation of the STN at $t = 100$ ms pushes the system from a low state to a sustained high state. Excitation of the GP at $t = 400$ ms returns the system to the low state (such excitation may arise from GP afferents or from internal cell properties; see Gillies & Willshaw 1998). (b) Increased inhibition of the GP at $t = 200$ ms increases the mean activities of both the STN and GP. (c) Oscillations. An increase in the inhibition (I^{STR}) at $t = 200$ ms results in system oscillation. (d) Increased excitation to the STN increases the mean activities of both the STN and GP. Parameters in all simulations: $\sigma_{\text{max}} = 500$ Hz, $\kappa = 0.3$ mV $^{-1}$, $x_{\text{th}} = 15$ mV, $\xi_{\text{max}} = 100$ Hz, $\eta = 0.2$ mV $^{-1}$, $y_{\text{th}} = 10$ mV, $\tau^{\text{STN}} = 6$ ms, $\tau^{\text{GP}} = 14$ ms, $a = 50$ $\mu\text{V Hz}^{-1}$, $b = 100$ $\mu\text{V Hz}^{-1}$, $d = 80$ $\mu\text{V Hz}^{-1}$, $c = 120$ $\mu\text{V Hz}^{-1}$. Variations: (a) $b = 140$ $\mu\text{V Hz}^{-1}$, $d = 40$ $\mu\text{V Hz}^{-1}$, $c = 10$ $\mu\text{V Hz}^{-1}$, (b) $I^{\text{CTX}} = 8$ mV. (c) $I^{\text{CTX}} = 9$ mV, $a = 54$ $\mu\text{V Hz}^{-1}$. (d) $I^{\text{STR}} = 0$ mV.

oscillations. An important factor is the value of a that determines the interconnectivity within the STN.

The system can exist in three different types of state: single stable, bistable and oscillatory states. We first give an informal explanation of these states. The STN and GP are subject to excitatory and inhibitory influences, mediated by external inputs and the connectivity parameters a , b , c and d . We consider the effects of gradually increasing the value of a , determining the amount of interconnectivity within the STN. When a is small, giving little or no self-excitation of the STN, inhibition predominates, keeping the firing rates of both nuclei low and stable; there is a single stable state. For a higher value of a , the stable states of the system reflect the balancing out of the excit-

atory and inhibitory influences on the two nuclei. Stable states emerge at both high and low firing rates. For example, when the system is at rest with low firing rates, short-term excitation of the STN from the cortex will increase the activity of the STN, pushing the entire system into a stable state with sustained high firing rates. Conversely, increased excitation of the GP whilst in this high state will increase the inhibition of itself and the STN, ultimately driving both back to their low states (figure 2a). Both the high- and low-activity states are stable, and separating them is a barrier which can be crossed by external drive. This bistability causes the system to act in a ‘switch-like’ manner. The oscillatory state could be regarded as alternating jumps, around the barrier, and between the

Figure 3. Phase plots of the nullclines for (x, y) showing the lines of $dx/dt = 0$ (red) and $dy/dt = 0$ (green) in the (x, y) plane, (a) under parameter conditions with no connectivity within the STN, i.e. $a = 0$ (all other parameters are the same as the bistable condition given in figure 2a), (b) under bistable parameter conditions, and (c) under parameter conditions leading to oscillations (see figure 2c). Stable points and limit cycles are indicated in grey.

high- and low-firing states that are not self-sustaining. The low frequency bursting-like oscillations emerge in both nuclei. Finally, for very large a or inputs I^{CTX} , excitation

Figure 4. Diagram illustrating the influence of cortical (I^{CTX}) and striatal (I^{STR}) inputs on oscillatory behaviour. The value of the fixed point $\hat{x}$ is plotted against I^{CTX} and I^{STR} . Blue indicates regions of stable fixed point solutions and green indicates regions of unstable fixed point solutions associated with oscillations, i.e. every point in the green region is a solution where the system is in stable oscillation. As inhibitory input from the striatum to the GP is increased, less excitatory cortical input to the STN is required to generate oscillations. Cortical input is plotted unconventionally on the z -axis and STN activity ($\hat{x}$) on the x -axis to highlight this relationship.

from the STN predominates to keep both nuclei firing at a high rate.

We now give a more mathematically based description of the system states and their properties.

(i) Interconnectivity within the STN

If there is no interconnectivity, the STN–GP complex simply responds to its external input by summing the afferent influences (from cortex and striatum); without interconnectivity it simply follows its transient input. The system resides at equilibrium and when disturbed will always return to this equilibrium. In this case A is positive since $a = 0$ and $A \equiv (1 - a\sigma'(x_0))/\tau^{\text{STN}}$. This is the condition as described in Appendix A § c(i), where the state (x_0, y_0) is stable (see phase plot figure 3a).

(ii) Influences of cortical and striatal input

In the presence of low levels of interconnectivity within the STN (small a) cortical excitation or striatal inhibition can lead to an increase in the equilibrium firing rates of the STN and GP. Paradoxically, increased inhibition of the GP can lead to an increase rather than a decrease in its mean firing rate (figure 2b). In particular, if a is sufficiently small that A is negative yet $A + B$ and $AB + CD$ both remain positive, all equilibrium points are stable (see § c(i) in Appendix A and equations (A 4) and (A 5)).

(iii) Conditions under which bistable behaviour occurs

In previous work (Gillies & Willshaw 1998) it was shown in a detailed model of the STN that it can display bistable behaviour. It can maintain activity at a high-firing rate or at a low-firing rate and can be driven between the two states by transient external input. A similar situation can occur in this coupled system. If a is sufficiently large

so that A is sufficiently negative to make $AB + CD$ negative (§ c(iii) of Appendix A), the state (x_0, y_0) becomes an unstable saddle point. Two new stable states then appear, which will attract any deviation from (x_0, y_0) (see phase plot figure 3*b* and figure 2*a*), one at a high-firing rate, the other at a low-firing rate.

(iv) *Conditions for the generation of oscillations*

With the system in the bistable state, a sufficient increase in the value of a causes two (one stable and one unstable) of the three equilibrium points to disappear, the remaining equilibrium point becoming unstable (§ c(ii) in Appendix A). As the values of (x, y) are bounded, a limit cycle around this unstable point results, giving rise to oscillations in the values of x and y (figures 2*c* and 3*c*). The mathematics relating to this condition are summarized in Appendix B. An important result is the relationship between external inputs and the onset of oscillations. As striatal inhibition is increased, *less* cortical input is needed for the system to enter the oscillatory state (figure 4).

(v) *Frequency of oscillations*

The frequency is proportional to $\sqrt{CD - (A - B)^2/4}$. This suggests that an increased x_0 (e.g. by increased cortical input) is likely to lead to lower oscillation frequencies due to a more negative $A < 0$. This is consistent with the converse observation that cutting the cortical input to the STN results in higher oscillation frequencies (Plenz & Kitai 1999).

4. DISCUSSION

In our model of the STN–GP interaction, three types of behaviour have been identified: a single stable state (which is the only possible state in the absence of interactions within the STN), a state in which the system acts as a bistable switch, being driven between the high and the low states by brief cortical or striatal stimulation, and a state comprising synchronized, low-frequency bursts of activity.

Assuming that the system has other functions than merely returning to a single stable state when disturbed from it, the bistable state could be interpreted as the normal mode of behaviour of the STN–GP system by means of which signals are provided from the STN to the two output nuclei of the basal ganglia to provide control, such as signals for terminating movement (Gillies & Willshaw 1998). The other mode of behaviour comprises synchronized, low-frequency bursts of activity. These could be identified with such movement disorders as Parkinsonism and could be interpreted either as one of the normally occurring pacemaker activities that are unmasked by Parkinsonian conditions (Plenz & Kitai 1999) or as an abnormally functioning state of the system.

Our model of the interactions between the GP and STN can be used to characterize the oscillations generated by such a system under these abnormal conditions and to isolate what parameters are important in their generation, as follows.

- (i) Assuming interactions within the STN, increased inhibition of the GP from the striatum is predicted to lead to an *increase* in firing rate in the GP. This

arises from the dynamic interaction between the STN and GP. Increased inhibition of the GP reduces its influence over the STN, allowing the interactions within the STN to increase STN activity resulting in increased activity in the GP (figure 2*b*). Increased cortical excitation of the STN also increases both the STN and GP firing rates (figure 2*d*); see § 3*a*(ii).

- (ii) The amount of excitatory cortical input to the STN required to generate oscillations depends on the amount of striatal input to the GP. As inhibitory striatal input to the GP is increased, the amount of excitatory drive required to generate oscillatory activity decreases (figures 4 and 2*c*; see § 3*a*(iv)). This prediction has important consequences for interpreting disorders of the basal ganglia. Under a current conceptual model of Parkinson's disease, striatal input to the GP is enhanced (Albin *et al.* 1989). Our analysis demonstrates how this can lead to increased oscillatory activity in the STN–GP complex. In animal models of Parkinson's disease, significantly increased oscillatory activity is observed in the STN and GP (Bergman *et al.* 1994; Nini *et al.* 1995). In human patients there is a significant population of rhythmic STN cells (Levy *et al.* 2000; Magariños-Ascone *et al.* 2000). This is a small population in comparison with observed phasic or tonic STN cell firing patterns (Magariños-Ascone *et al.* 2000). Our model demonstrates the circuits that may underlie the generation of the rhythmic patterns while not ruling out other circuits and inputs that will allow cells to deviate from this rhythmic activity (e.g. via direct cortical input to the STN).
- (iii) The frequency of repetition of the high-frequency bursts can vary in the model and it depends upon the strengths of connections between structures, time constants of the neurons and the external inputs (see § 3*a*(v)). The range of oscillation frequencies observed within biologically realistic parameter regions in computer simulation is 3–25 Hz. This fits well with burst frequencies found in *in vivo* monkey recordings under Parkinsonism-type conditions (Nini *et al.* 1995) and in recordings from human patients (Levy *et al.* 2000; Magariños-Ascone *et al.* 2000). This differs from cell culture studies (Plenz & Kitai 1999) where the strength and extent of connections and external inputs probably differ significantly from those found in *in vivo* recordings.

APPENDIX A: MATHEMATICAL DETAILS

(a) *Formulation*

The coupled differential equations describing the membrane potentials x and y of the STN and GP neurons specify the ensemble average of a spike train model where impulses are convolved with an alpha-type response function (see Shamma 1989). This allows us to model spike trains in a dynamic manner as follows:

$$\begin{aligned}\tau^{\text{STN}} dx/dt &= -x + a\sigma(x) - c\xi(y) + I^{\text{CTX}} \\ \tau^{\text{GP}} dy/dt &= -y - b\xi(y) + d\sigma(x) + I^{\text{STR}},\end{aligned}\tag{A 1}$$

where a , b , c and d specify the couplings within and between STN and GP; I^{CTX} and I^{STR} represent cortical and striatal influences. The STN and GP potentials x and y are related to firing frequency $\sigma(x)$, $\xi(y)$ via sigmoidal functions:

$$\sigma(x) = \frac{\sigma_{\max}}{1 + e^{-\kappa(x - x_{\text{th}})}}, \quad \xi(y) = \frac{\xi_{\max}}{1 + e^{-\eta(y - y_{\text{th}})}}. \quad (\text{A } 2)$$

(b) *Equilibrium points*

The equilibrium values of (x, y) in equation (A 1) are found by setting $dx/dt = dy/dt = 0$. The nature of these equilibria can be obtained by looking at the linear deviations $(\hat{x}, \hat{y})$ of (x, y) from the equilibrium point (x_0, y_0) as expressed by the matrix equation

$$\begin{pmatrix} d\hat{x}/dt \\ d\hat{y}/dt \end{pmatrix} = M \begin{pmatrix} \hat{x} \\ \hat{y} \end{pmatrix},$$

where

$$M = \begin{pmatrix} -A & -C \\ D & -B \end{pmatrix},$$

$$A \equiv (1 - a\sigma'(x_0))/\tau^{\text{STN}},$$

$$B \equiv (1 + b\xi'(y_0))/\tau^{\text{GP}},$$

$$C \equiv c\xi'(y_0)/\tau^{\text{STN}},$$

$$D \equiv d\sigma'(x_0)/\tau^{\text{GP}},$$

and $\sigma'(x_0)$ and $\xi'(y_0)$ are the derivatives of σ and ξ at (x_0, y_0) . The deviations follow the solution $\hat{x} \propto e^{\lambda_{1,2}t}$, where $\lambda_{1,2}$ are the eigenvalues of M

$$\lambda_{1,2} = -\frac{1}{2}(A + B) \pm \frac{1}{2}\sqrt{(A + B)^2 - 4(AB + CD)}. \quad (\text{A } 3)$$

(c) *Nature of the equilibrium points*

It is useful to distinguish the following three cases, which differ with respect to the signs of $A + B$ and $AB + CD$. Note that B , C and D are restricted to positive or zero values, whereas A , which depends on a , the amount of interconnectivity within STN, can have both positive and negative values.

(i) $A + B > 0$ and $AB + CD > 0$. The eigenvalues $\lambda_{1,2}$ are real or complex with negative real parts, giving a stable equilibrium point.

(ii) $A + B < 0$ and $AB + CD > 0$. Both eigenvalues have positive real parts, giving an unstable equilibrium point. When the system has a single equilibrium point that is of this type, limit cycles (oscillations) result, as specified by the Poincaré–Bendixson theorem (Edelstein-Keshet 1988).

(iii) $AB + CD < 0$. Both eigenvalues are real but with opposite signs, yielding an unstable saddle point from which deviations will reach one of the two separate basins.

(d) *Effect of changes in external inputs*

Changing the cortical and striatal inputs by δI^{CTX} and δI^{STR} changes the equilibrium potentials (x_0, y_0) by

$$\delta x_0 = \frac{B(\delta I^{\text{CTX}}/\tau^{\text{STN}}) - C(\delta I^{\text{STR}}/\tau^{\text{GP}})}{AB + CD}, \quad (\text{A } 4)$$

$$\delta y_0 = \frac{D(\delta I^{\text{CTX}}/\tau^{\text{STN}}) + A(\delta I^{\text{STR}}/\tau^{\text{GP}})}{AB + CD}. \quad (\text{A } 5)$$

APPENDIX B: MATHEMATICS OF OSCILLATIONS

(a) *Existence of oscillations*

Appendix A shows that oscillation occurs around a particular type of unstable fixed point that has $A + B < 0$ and $AB + CD > 0$. This unstable fixed point can arise out of a stable fixed point only by decreasing A to a negative value (since $B, C, D > 0$). This can be achieved by increasing interconnection a within the STN or increasing $\sigma'(x_0)$ by shifting x_0 via external drives I^{CTX} and I^{STR} . In the case in figure 2c, the parent (stable) fixed point is the one in figure 2b, and is the only fixed point of the system (figure 3a). In other cases, the parent fixed point can be one of the two stable fixed points among the three fixed points in the bistable condition (figure 3b). For instance, in our simulations, increasing a causes the unstable equilibrium point to coalesce with the lower stable fixed point. Increasing a further causes both to disappear, leaving just one equilibrium. This happens at the point at which the two curves $dx/dt = 0$, $dy/dt = 0$ drawn in the (x, y) plane have identical gradients near the fixed points. In both of these cases, the single remaining equilibrium point (with $AB + CD > 0$) is unstable, taking complex eigenvalues with positive real values (§ c(ii) in Appendix A). According to the Poincaré–Bendixson theorem (Edelstein-Keshet 1988), this implies the existence of limit cycles if both x and y are bounded. This is true in this case as both (x, y) are bounded from above and below. For example, from equation (A 1), x is bounded as the sigmoidal functions are bounded $0 < \sigma(x) < \sigma_{\max}$.

(b) *The onset of oscillations*

As a is increased, which causes a decrease in A , the value of $A + B$ for the stable fixed point moves from positive to zero, at which point oscillations occur:

$$A + B = (1 - a\sigma'(x_0))/\tau^{\text{STN}} + (1 + b\xi'(y_0))/\tau^{\text{GP}} = 0. \quad (\text{B } 1)$$

Under STN and GP resting conditions (rest firing rates *ca.* 15 and *ca.* 32 Hz, respectively (Féger *et al.* 1991; Kita & Kitai 1991)), x_0 is at the base of its sigmoid function $\sigma(x_0)$ giving a small value of $\sigma'(x_0)$. Increasing (x_0, y_0) by stronger cortical input I^{CTX} or striatal inhibition ($-I^{\text{STR}}$) leads to increased $\sigma'(x_0)$ eventually to give $A + B = 0$ (equation (B 1)). This is demonstrated in the bifurcation diagram of figure 4. Either increased cortical input to the STN or increased striatal input to the GP can help bring about oscillations.

(c) *Frequency of oscillations*

The frequency of the oscillations is approximated by the imaginary part of the eigenvalues $\lambda_{1,2}$. It is proportional to $\sqrt{CD - (A - B)^2/4}$.

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